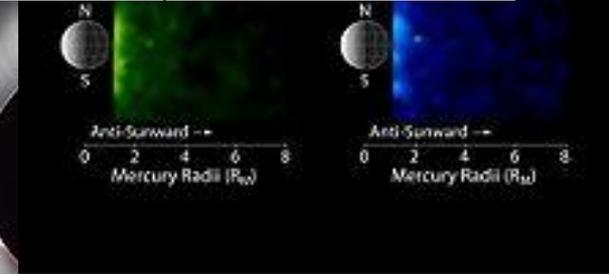
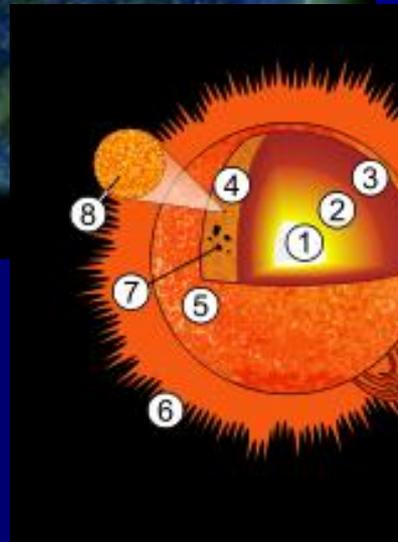
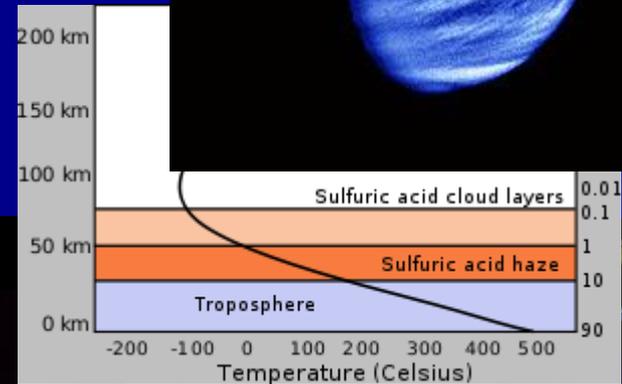
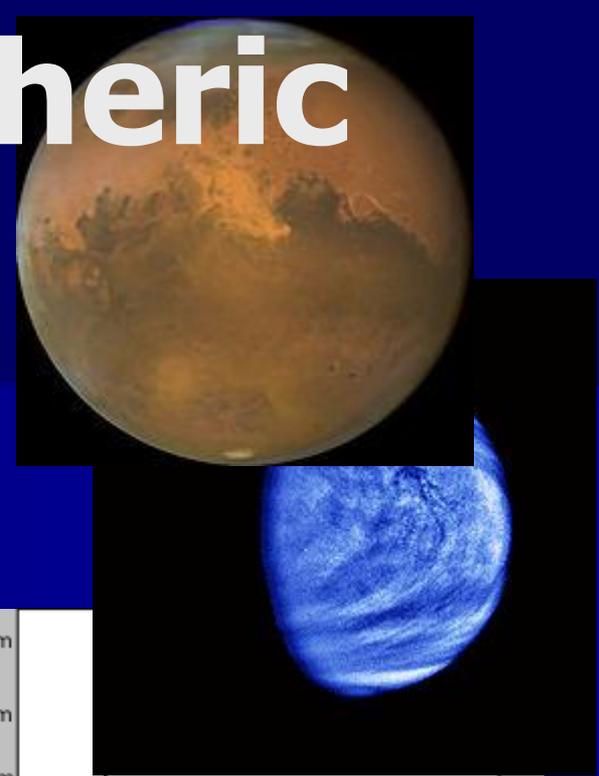
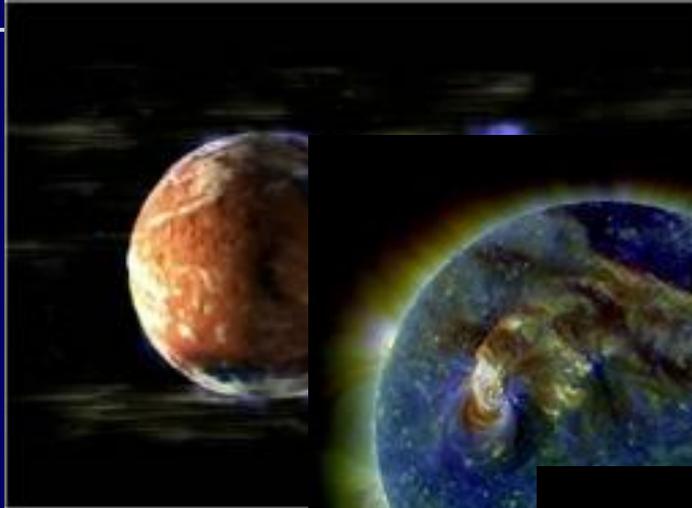


Space Physics

Space & Atmospheric Physics



Lecture – 13 A

Radio Wave Communication

Radio waves

Radio Communication

Reflection of Radio Waves

Absorption of Radio Waves

Complex Refractive Index

Reflection Heights

Deviating Region Absorption, Non-

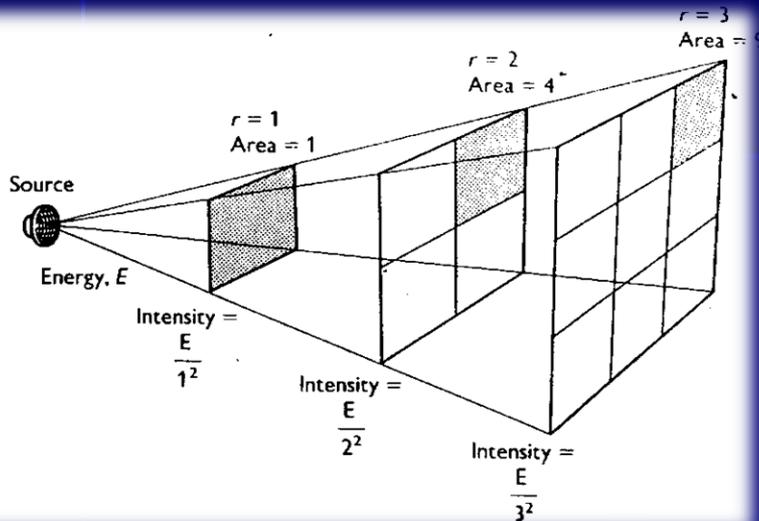
Deviating Region Absorption

Ordinary/Extra Ordinary Waves

Ionosphere – Sounding Techniques

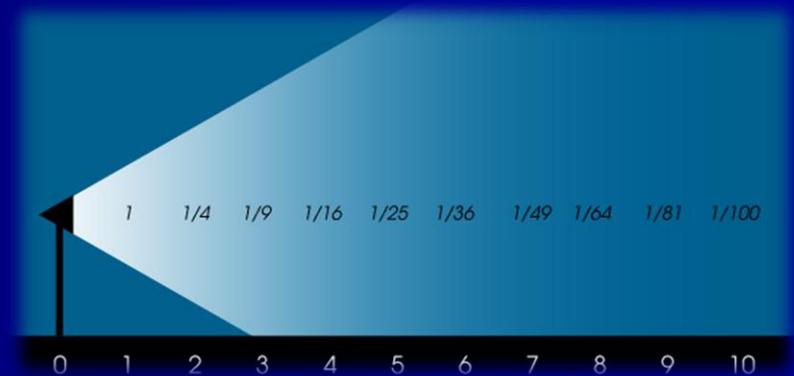
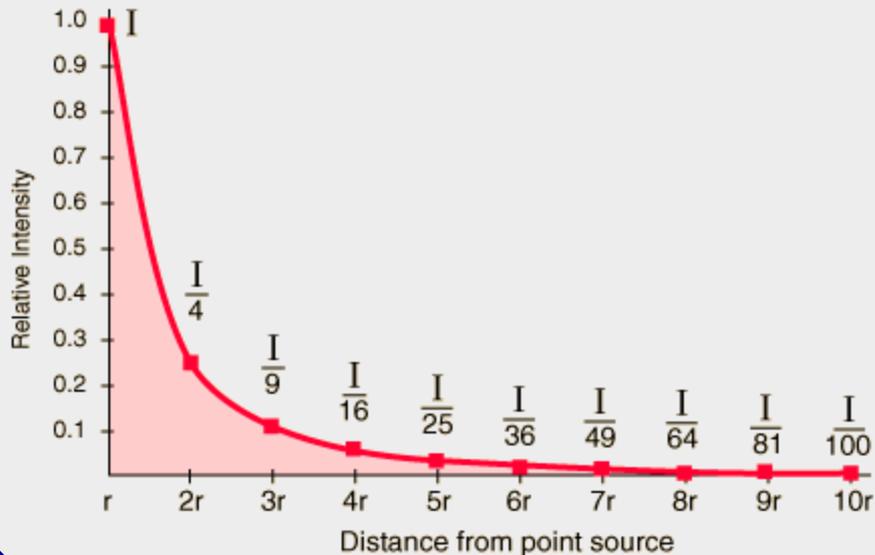
Pulse Reflection Methods

Free Space Propagation



In free space all EM waves obey the **Inverse Square law** which states that the **power density of an EM wave** is proportional to the **inverse of the square of the distance from the Source**.

$$\rho_p \propto \frac{1}{r^2}$$

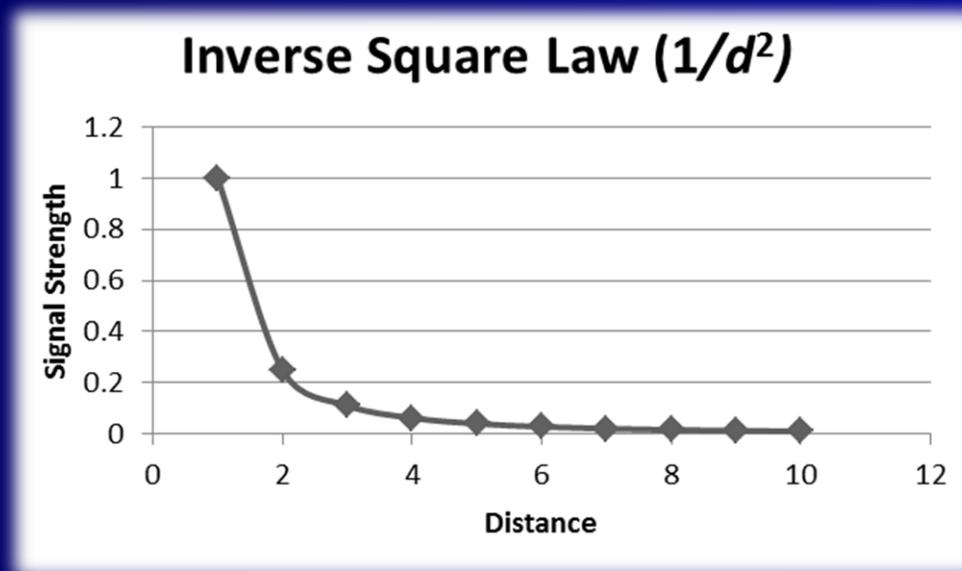


Distance from point source

Free Space Propagation

Doubling the distance from a transmitter means that the power density of the radiated wave at that new location is reduced to one quarter of its previous value !

The power density per surface unit is proportional to the product of the electric and magnetic field strengths. Thus, **doubling the propagation path distance from the transmitter reduces each of their received field strengths over a free-space path by one half.**



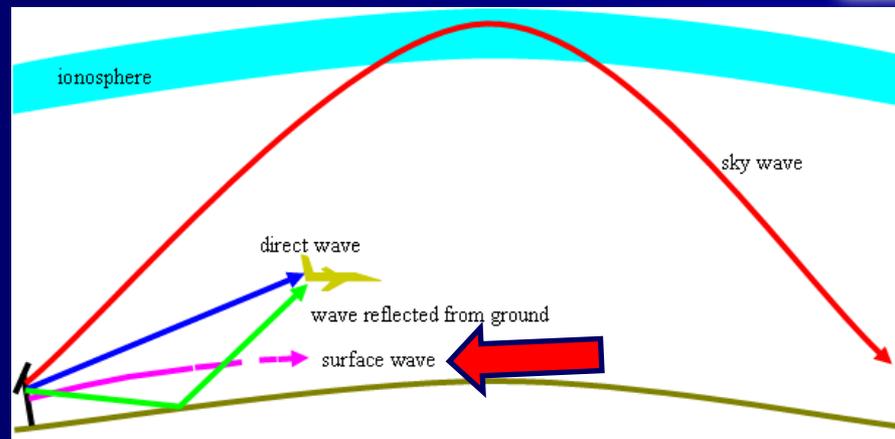
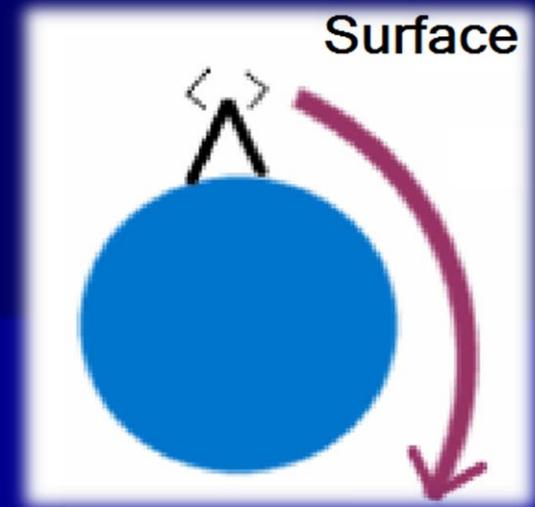
Modes

Band		Frequency	Wavelength	Propagation via
VLF	Very Low Frequency	3–30 kHz	100–10 km	Guided between the earth and the ionosphere.
LF	Low Frequency	30–300 kHz	10–1 km	Guided between the earth and the D layer of the ionosphere. Surface Waves.
MF	Medium Frequency	300–3000 kHz	1000–100 m	Surface waves. E, F layer ionospheric refraction at night, when D layer absorption weakens.
HF	High Frequency (Short Wave)	3–30 MHz	100–10 m	E layer ionospheric refraction F1, F2 layer ionospheric refraction
VHF	Very High Frequency	30–300 MHz	10–1 m	Infrequent E ionospheric refraction. Extremely rare F1, F2 layer ionospheric refraction during high sunspot activity up to 80 MHz. Generally direct wave. Sometimes tropospheric ducting
UHF	Ultra High Frequency	300–3000 MHz	100–10 cm	
SHF	Super High Frequency	3–30 GHz	10–1 cm	
EHF	Extremely High Frequency	30–300 GHz	10–1 mm	Direct wave limited by absorption.

Propagation Modes

Surface Modes :

Lower frequencies (between 30 kHz – 3000 kHz) have the property of the curvature of the Earth via ground wave propagation in the majority of occurrences.

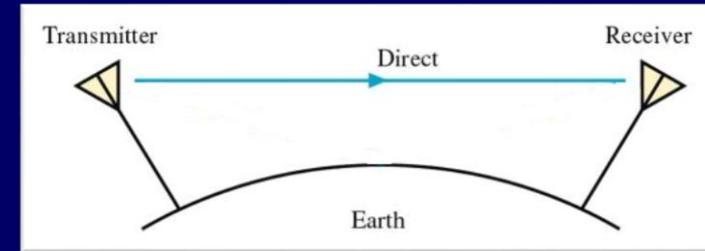


Ground waves are attenuated rapidly as they follow the Earth's surface. Attenuation is proportional to the frequency making this mode mainly useful for LF and VLF frequencies.

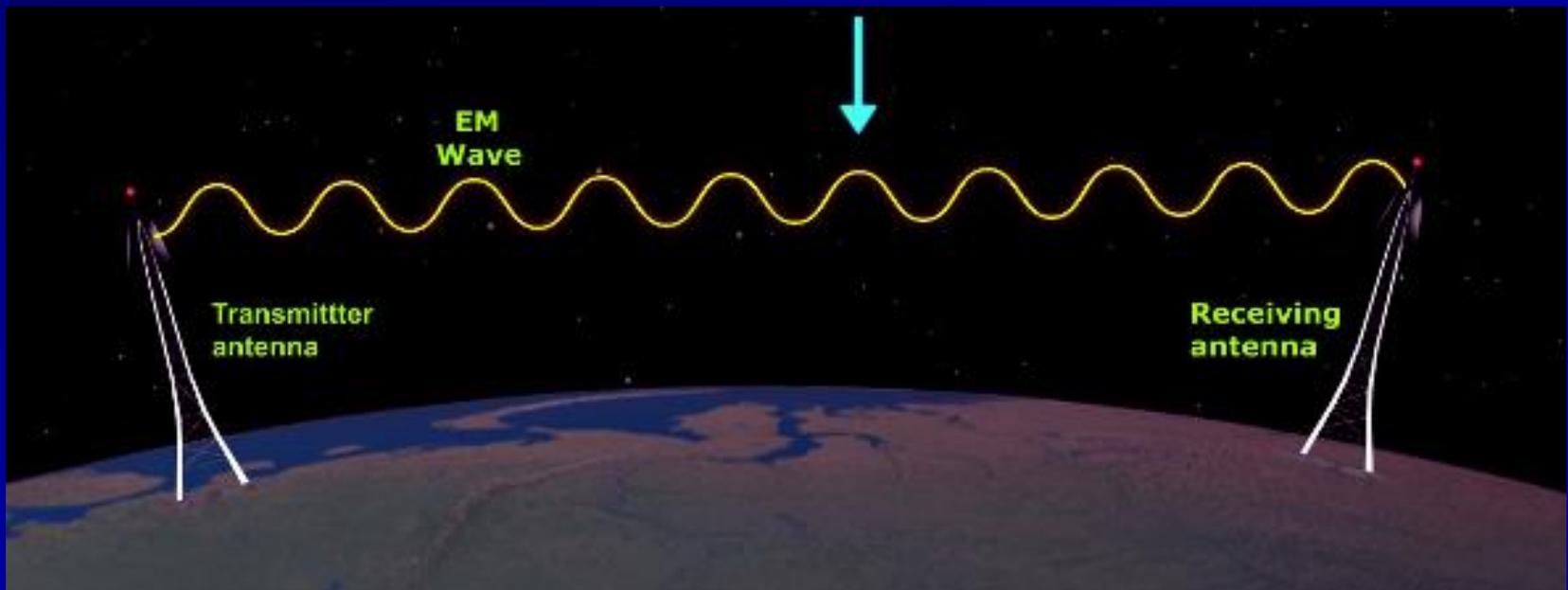
$$\text{Attenuation} \propto f$$

Propagation Modes

Direct Modes : (line-of-sight)



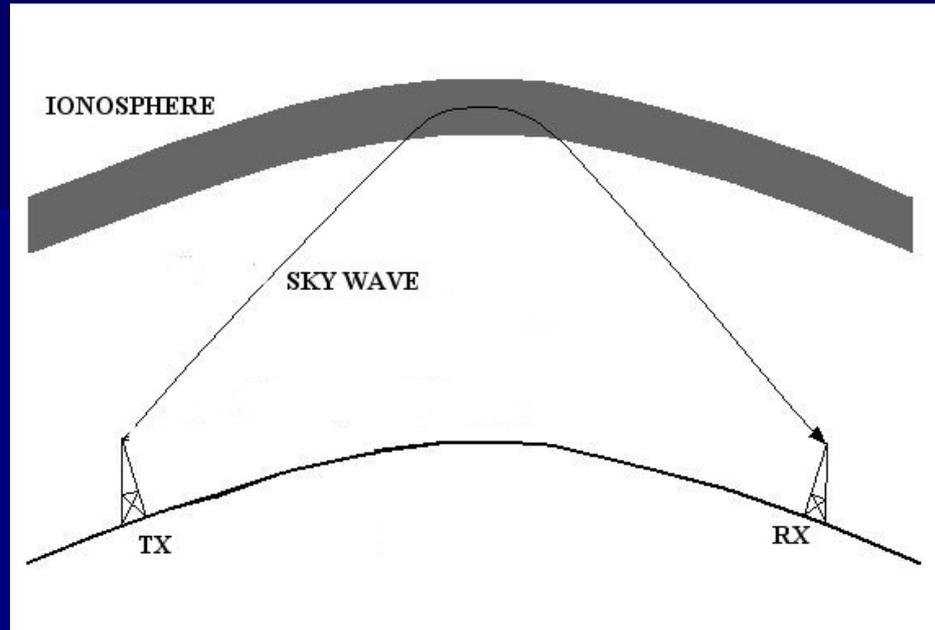
Line of sight is the **direct propagation of radio waves between antennas** that are visible to each other. This is probably the most common of the radio propagation modes **at VHF and higher frequencies**



Propagation Modes

Ionospheric Modes :

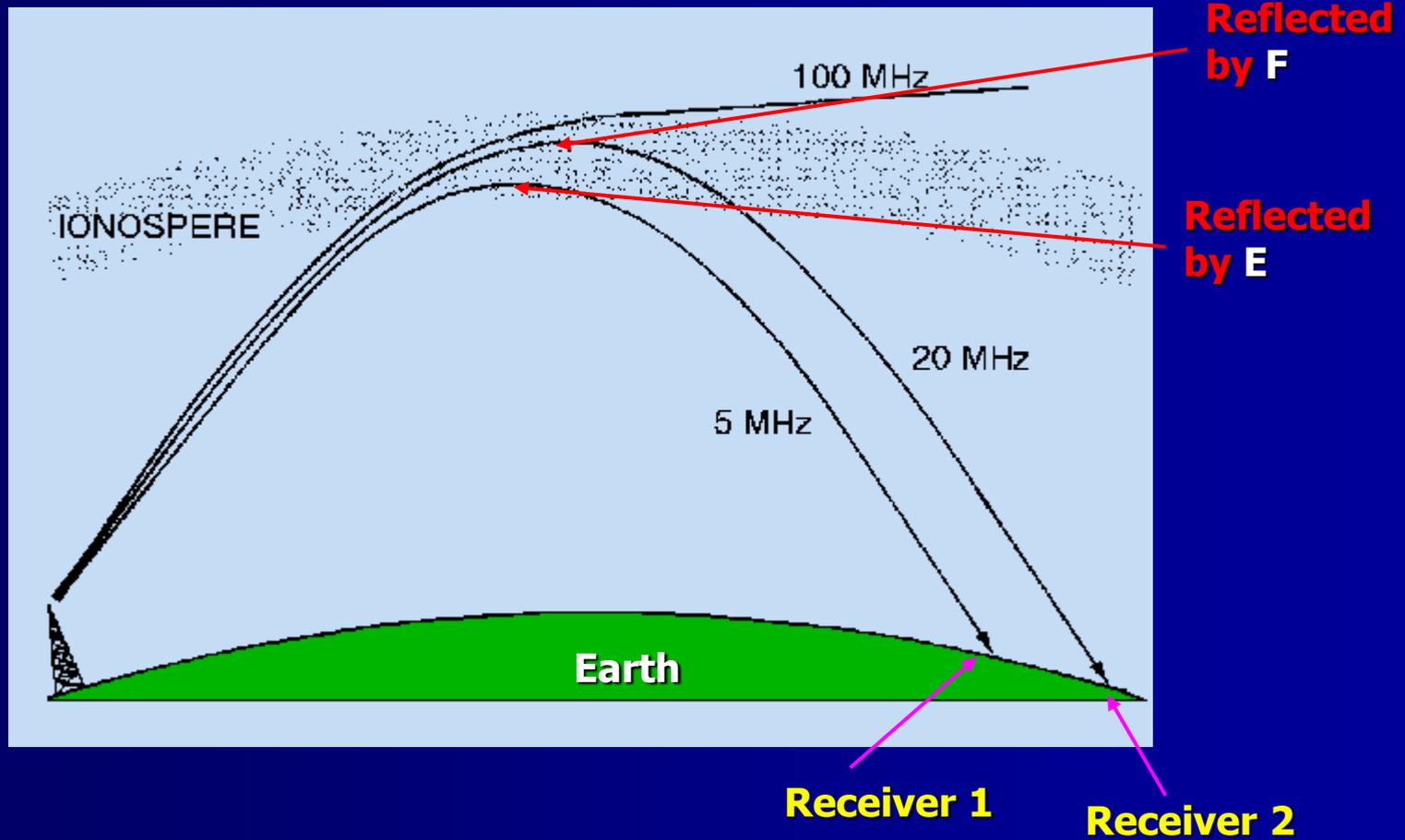
Sky wave propagation, also referred to as skip, is any of the modes that rely on **refraction** of radio waves in the **ionosphere**, which is made up of one or more ionized layers in the upper atmosphere



F₂ layer is the most important ionospheric layer for **HF** propagation, through **F₁**, **E** and **D** layers also play some role. These layers are directly affected by the Sun on a daily cycle, the reasons and the **11-year sunspot cycle** determines the utility of these modes. During **solar maxima**, the whole **HF range** up to **30 MHz** can be used and **F₂** propagation up to **50MHz** are observed frequently depending upon daily solar flux values. During **solar maxima**, **propagation of higher frequencies is generally worse**.

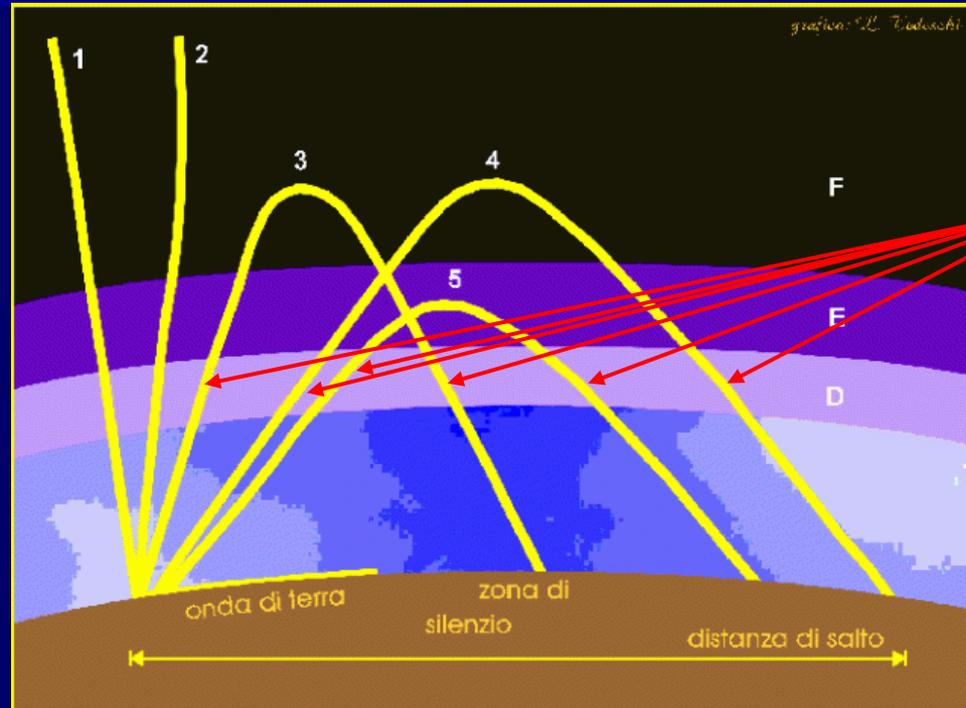
Reflection of Radio Waves

Generally, radio waves reflected by **E** and **F** regions, because that ionosphere E and F have more electron densities compared to the region **D**.



Reflection of Radio Waves

- Always radio waves **absorbed by the region D** in the ionosphere. That is a main problem of the radio communication.

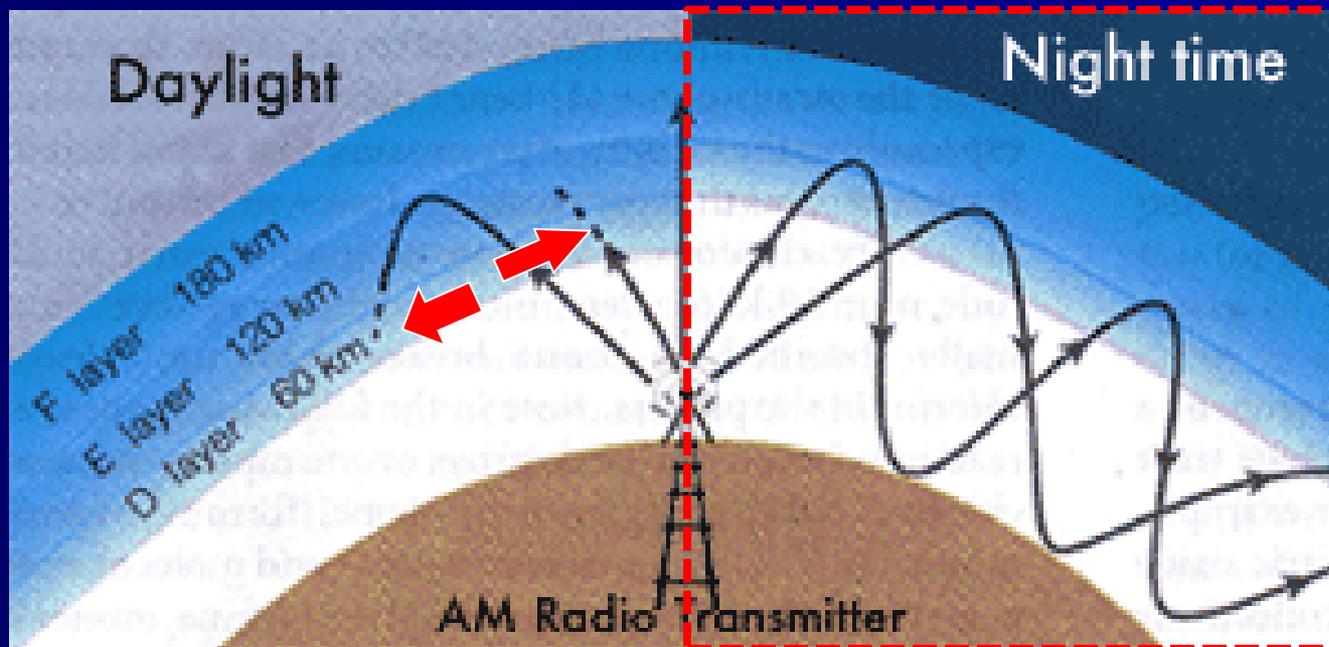


Absorbed
by **D**

Because in the day time, radio waves reflected from the region E and F. But always the waves are going through the region D twice. (incident beam and reflected beam)

Reflection of Radio Waves

- In the Night time region **D** of the ionosphere vanishing suddenly, because the Sun goes down. But **E** and **F** regions are not destroying suddenly, specially the region **F** (F₂)

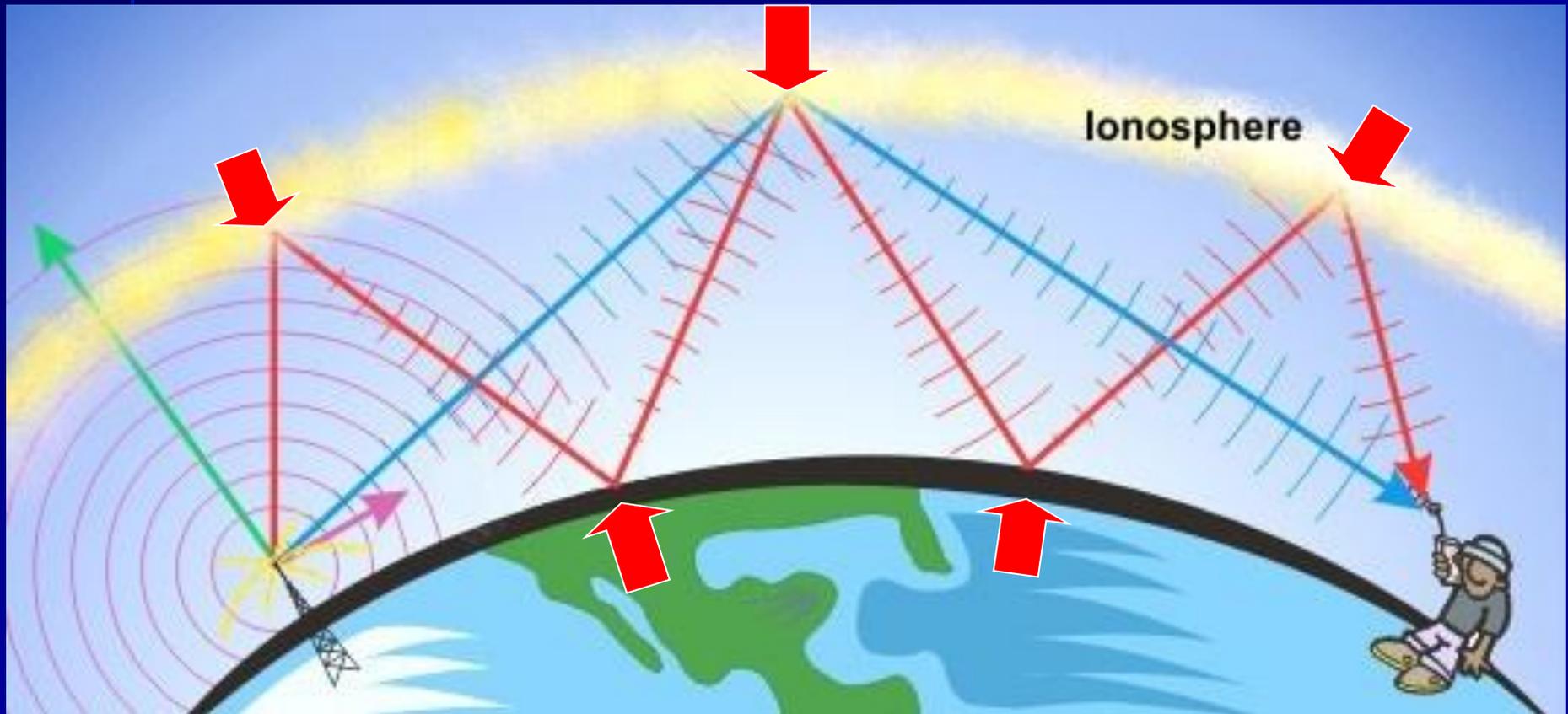
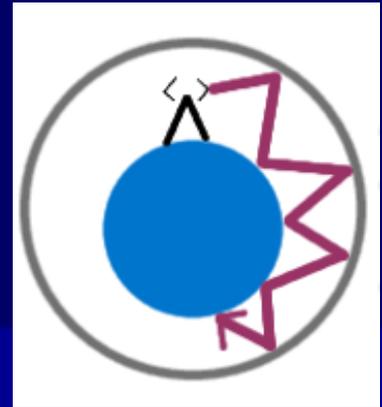


As a result SW radio broadcasting lines can listen very clearly at the night.

Eg: BBC Sandeshaya

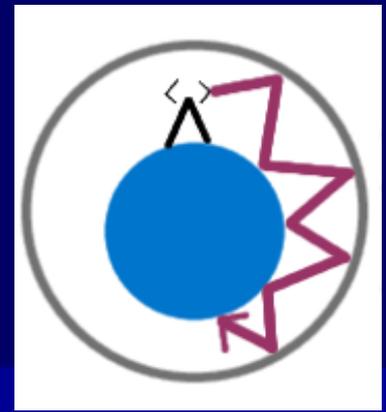
Multi - Reflection

- In the Night there are high possibilities for multi reflection too.



Multi - Reflection

Then the Sun is rising in the morning, all communications reflected from F-region going to weak. Because again D-region is creating !



In the night MW (medium waves) also reflected from the region F. As a result, we can listen some MW channels from India.

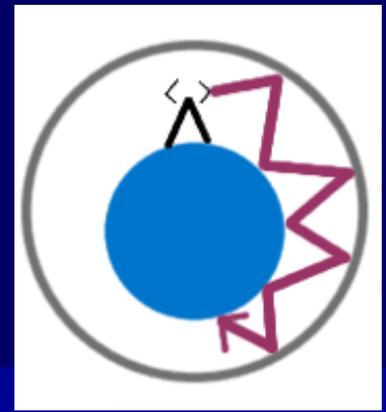
For long distance communication it is very useful SW reflection broadcast method. (But the problem is D-region in the day)

We can not use surface waves communication for long distance communications. There are two reasons:

01. The problem from the curvature of the earth.
02. The surface waves dies more quickly as the frequency increases;

$$\text{Range (km)} = \frac{200}{\sqrt{\text{Frequency (MHz)}}$$

Multi - Reflection



$$\text{Range (km)} = \frac{200}{\sqrt{\text{Frequency (MHz)}}$$

If **f** increases the range of the propagation, **R** decreases rapidly. But **f** decreases the range of the propagation, **R** increases.

Eg: A 100 kHz low frequency radio wave use for the surface propagation to long distance. Find the maximum range of the propagation of the wave.

$$f = 100 \text{ kHz} \longrightarrow f = 0.1 \text{ MHz}$$

Using, $R \text{ (km)} = \frac{200}{\sqrt{f \text{ (MHz)}}$ \longrightarrow $R \text{ (km)} = \frac{200}{\sqrt{0.1}}$

$$R = 632.456 \text{ km}$$

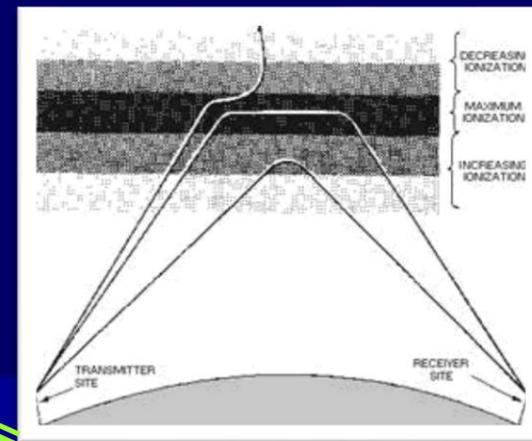
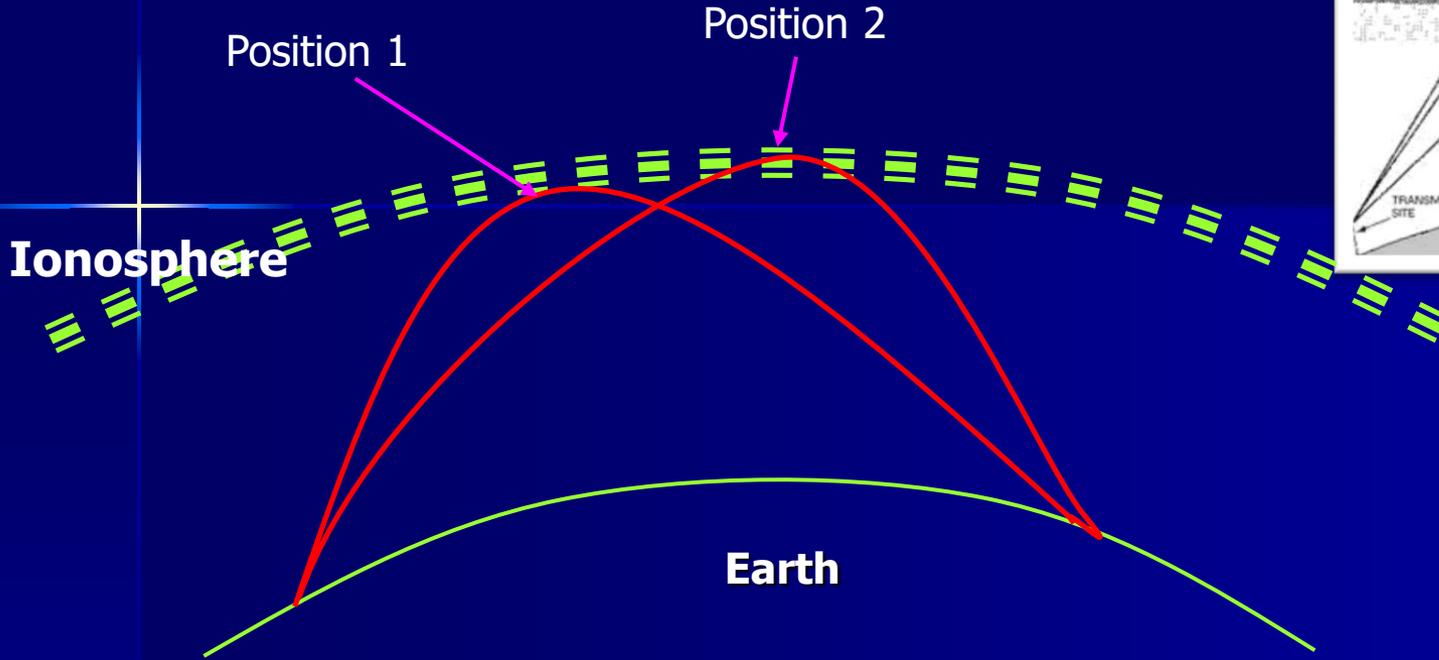
$$R \approx 600 \text{ km}$$

After that range it is dying more quickly !

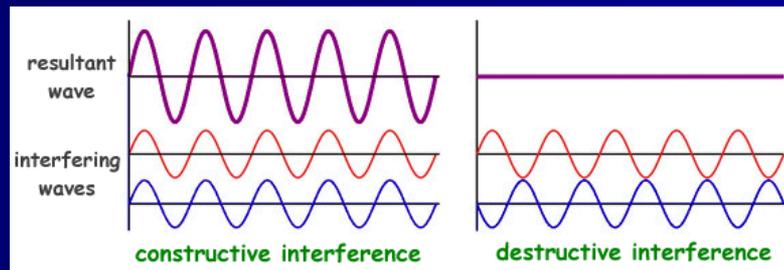
Properties of Reflection of Radio Waves

- Signal reflected from ionized layer of the ionosphere back down to Earth
- Reflection effect caused by refraction
- Reflection occurs when signal encounters a surface that is large relative to the wavelength of the signal.
- The radio wave is not reflected from a single point on the reflector but rather from an area on its surface. The size of the area required for reflection to take place depends reflecting substance.
- Radio wave may be reflected from various substances or objects they meet during travel between the transmitting and receiving sites.
- When the radio waves are reflected from flat surfaces, a phase shift in the alternations of the wave occurs.
- The shifting the phase relationship of reflected radio waves is one of the major reasons for fading.

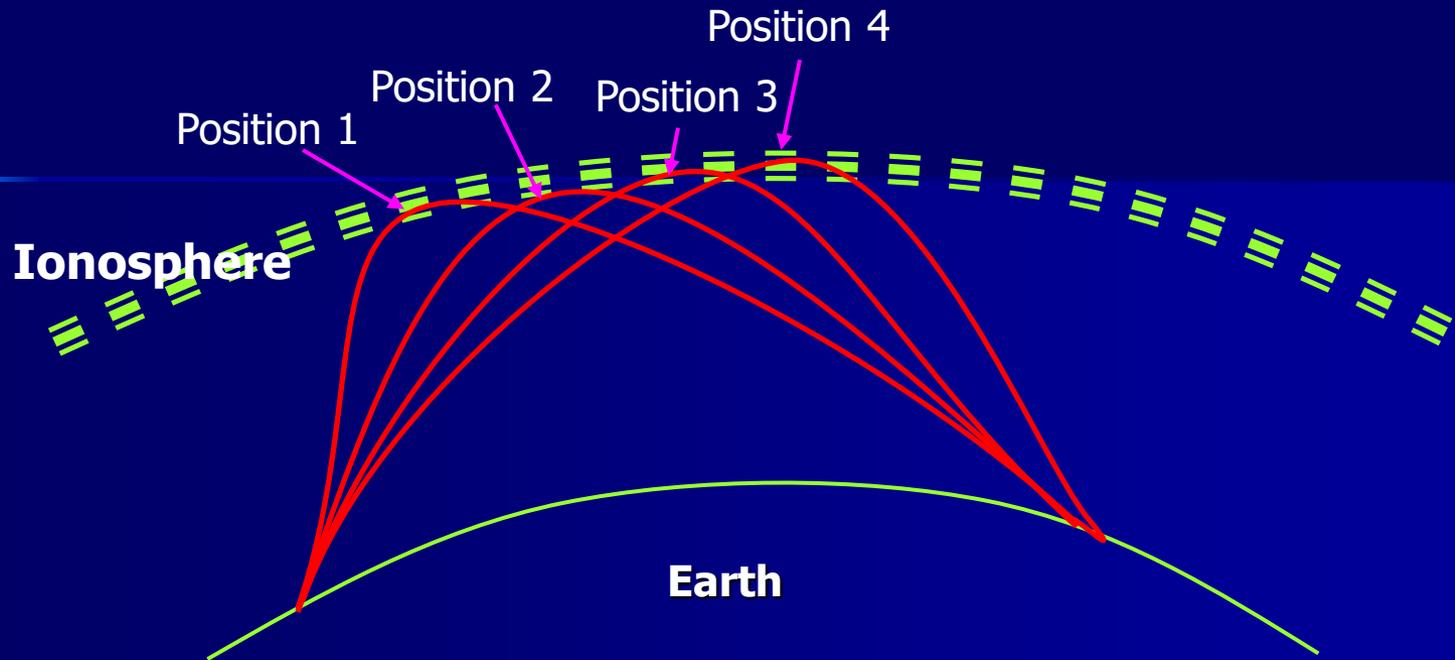
Multi – Position Reflection



This is a possible situation. That means multi position of the ionosphere is supported to reflected the radio wave to the same receiver. **As a result the strength of the signal may be very strong.** But sometimes it may be very weak, because the **result of interference** final wave may be cancelled.



Multi – Position Reflection

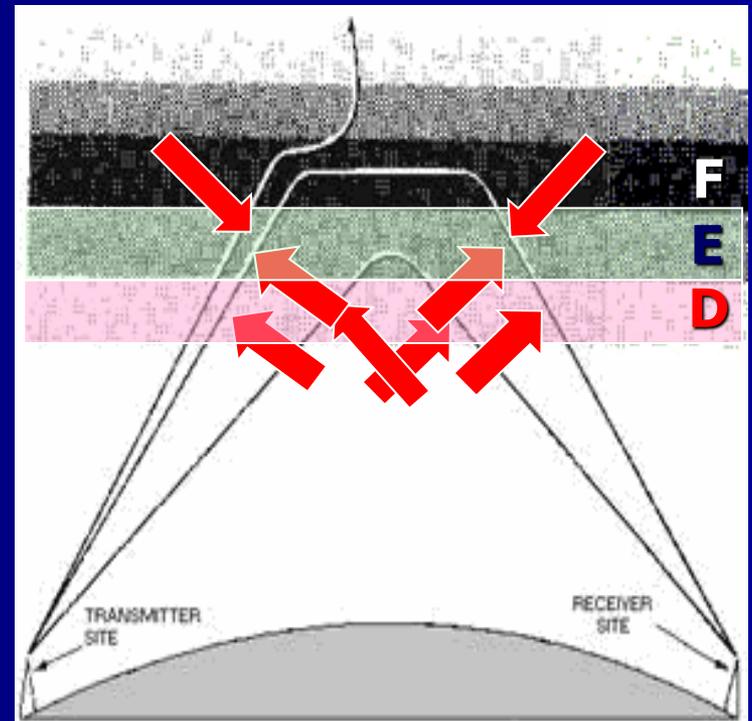


As a result of reflection of multi position final wave may be **destroyed**, because of the **attenuation of the carrier wave** of the radio signal. OR **final wave may be destroyed, because of anti-phase situation** (destructive interference)

Absorption of Radio Waves

Normally ions (electrons and positive ions) in the ionosphere interact with radio waves (EM waves). As a result, ions in the ionosphere are oscillating due to absorption energy from the radio waves. That means energy of the ions increases. Because, electric field of the radio wave transfer energy to ions. The temperature of the ionosphere is increases.

If we want to cancel this phenomena **we should use High Power radio waves** for this purpose. In the day time this phenomena is happening well, because, molecular density of the ionosphere is very high. Normally ions in the D-region are participated for this incident.



Absorption of Radio Waves

Oscillation frequency is the most important fact to absorb radio waves from the D-region. (**Oscillation frequency is no of oscillation per second of the electrons or positive ions**)

If **the radio wave frequency is equal to the oscillation frequency** of the D-region; the **wave is absorbed from the D-region**. Therefore, this **oscillation frequency is the most important factor for the Radio Wave Propagation**.

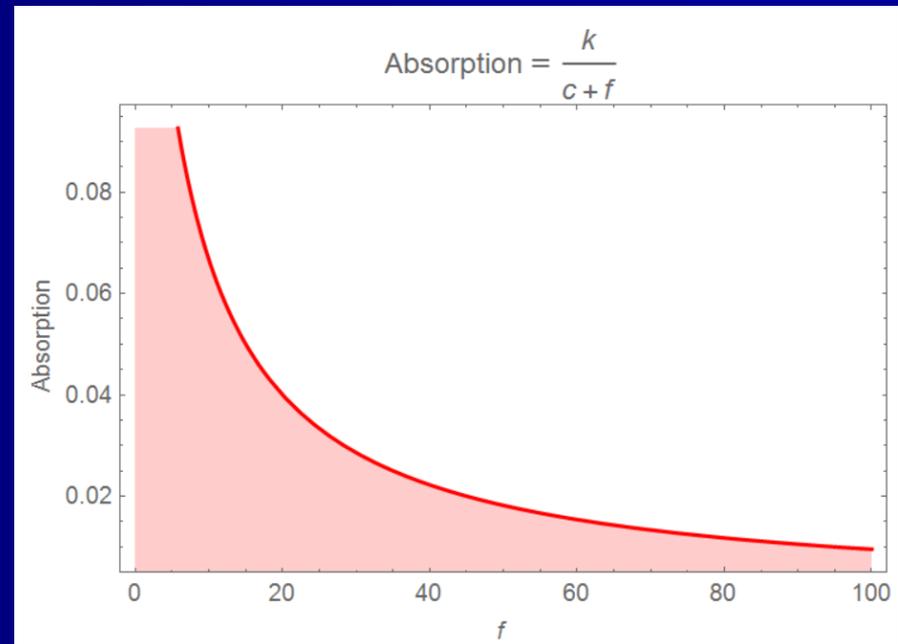
Normally, absorption is directly depend on the frequency of the Radio Wave.

$$\text{Absorption} \propto \frac{1}{c + \text{Frequency of the wave}}$$

Where, c is a constant.

$$Ab = \frac{k}{c + f}$$

Where, k is proportional constant.



Absorption of Radio Waves

$$\text{Absorption} \propto \frac{1}{c + \text{Frequency of the wave}}$$

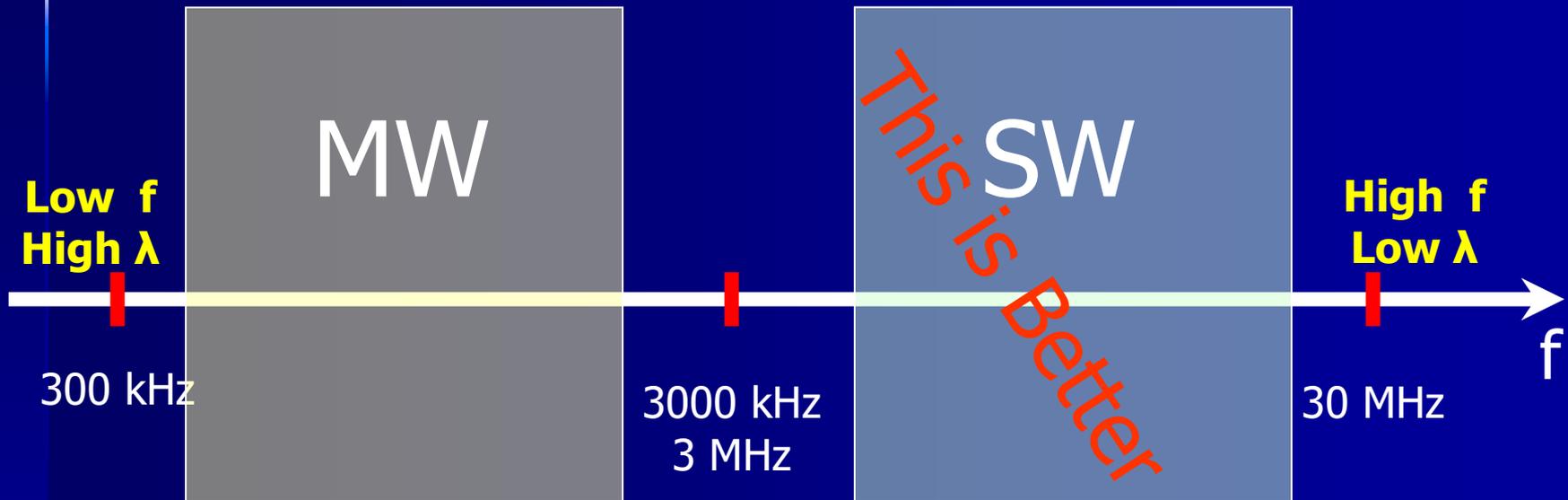
If **f** ↑ then **ab** ↓



If **f** ↓ (**λ** ↑) then **ab** ↑

Absorption of Radio Waves

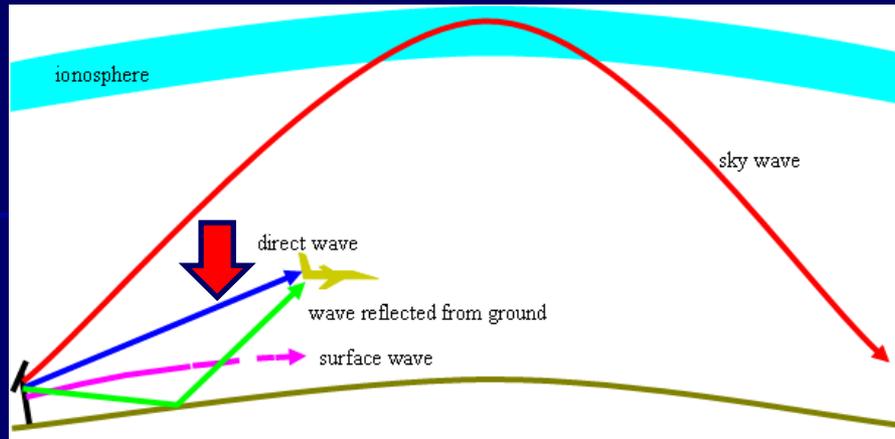
Generally, **Short Waves (SW)** are not **much absorbed** from the D-Region of the ionosphere. But, **Medium Waves (MW)** are **normally absorbed** from the ionosphere



In the day time, **MW** is almost absorbed from the D-Region of the ionosphere. But **SW** are not absorbed.

Therefore, **we can use MW to direct communication** without reflection from the ionosphere. That direct communication means use to **communicate by direct ground waves**.

Absorption of Radio Waves



But direct ground waves also face **some practical problems** due to **diffraction**. Because the wave is bending due to diffraction. **Therefore, we can not use ground waves to communicate too far.**

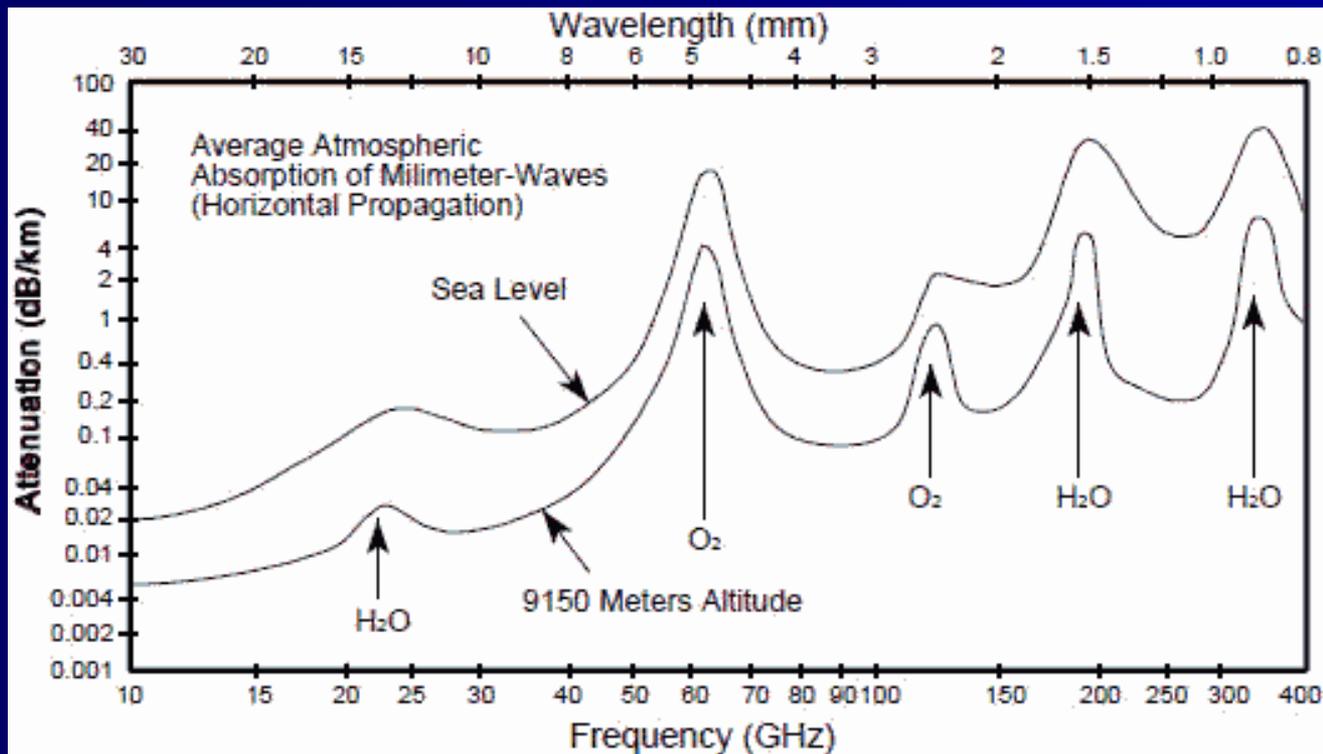
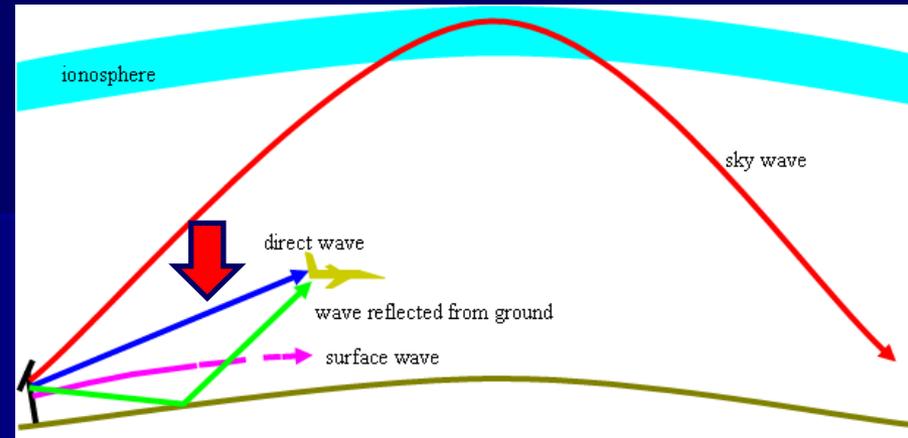
Frequency of the **MW** signal **very low** with comparative to others. As a result **range of propagating** is **high** for low frequency signals.

$$\therefore \text{Range (km)} = \frac{200}{\sqrt{f \text{ (MHz)}}}$$

But the problem is **diffraction of the wave.**

Absorption and attenuation of Radio Waves

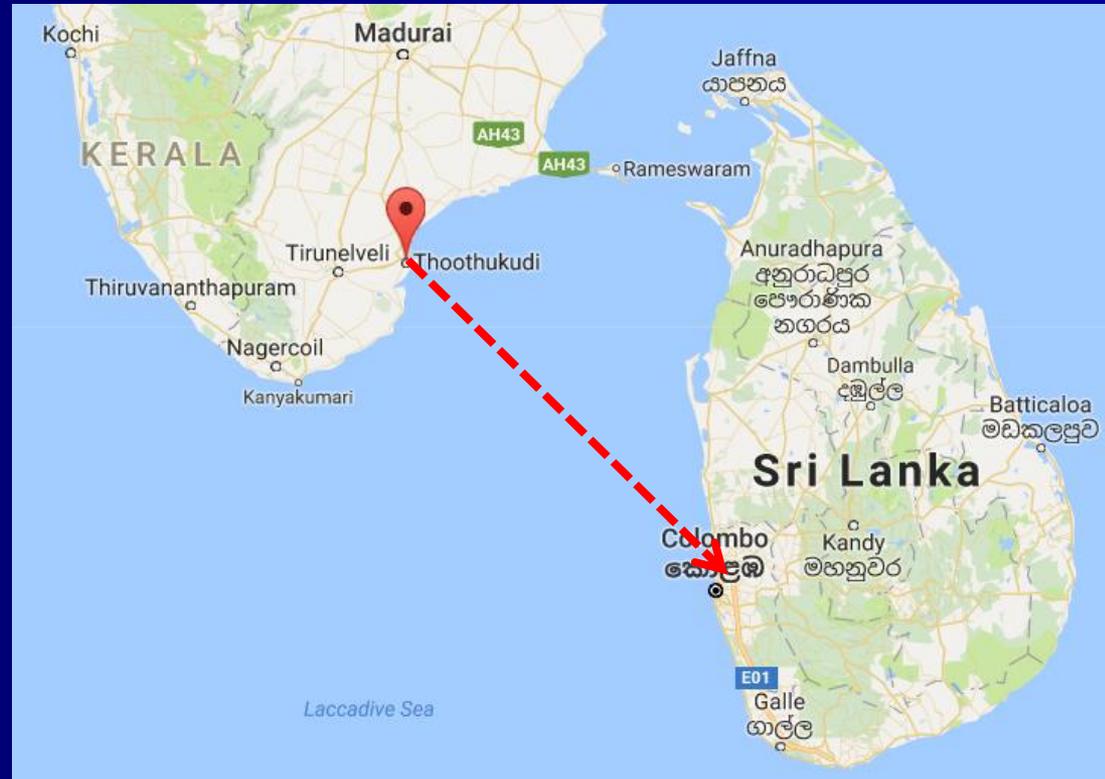
Direct ground waves face another problem due to many different gasses, water (water vapor) and particles contained the atmosphere. That is attenuation of the wave by atmosphere of the Earth.



Absorption and attenuation of Radio Waves

A significant atmosphere effect is that of **attenuation due to rain**. Below about **10 GHz**, rain fading is not very significant, but at **higher microwave frequencies, it becomes the major factor limiting path length**, particularly in areas that experience high level of rainfall (eg: **Watawala town in Sri Lanka**). In addition to attenuation of EM waves, rain and other precipitation (dew,...) tend to cause depolarization of the wave.

There is a ULF (Ultra Low Frequency) radio frequency channel from **Thoothukudi** in India whose frequency range is **1020 Hz – 1070 Hz**. We can listen to that channel in Sri Lanka because they are using **Very High Power Transmitter** !

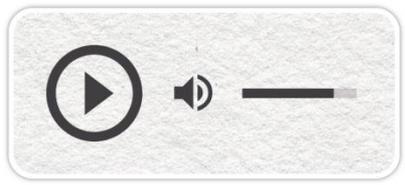




AIR Tuticorin

Language: **All India Radio / Tamil**

★★★★★ 5(1)



AIR Tuticorin (Thoothukudi) is part of All India Radio broadcasting live from Tuticorin(Thoothukudi) Tamilnadu broadcasting at 100.5 MHZ. All India Radio Tuticorin plays Akashvani news, Talk programmes, Current Affairs, World news and plays Tamil songs and Tamil talk programmes. Listen to Akashvani Tuticorin FM radio station online.

How useful was the programs in this Radio?

Select a star to rate it!



Average rating 5 / 5. Vote count: 1

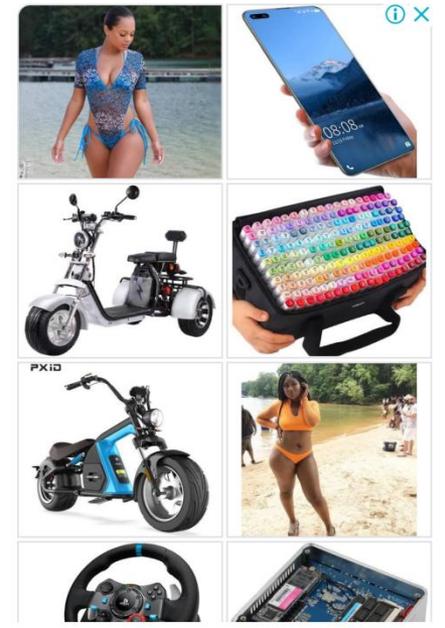
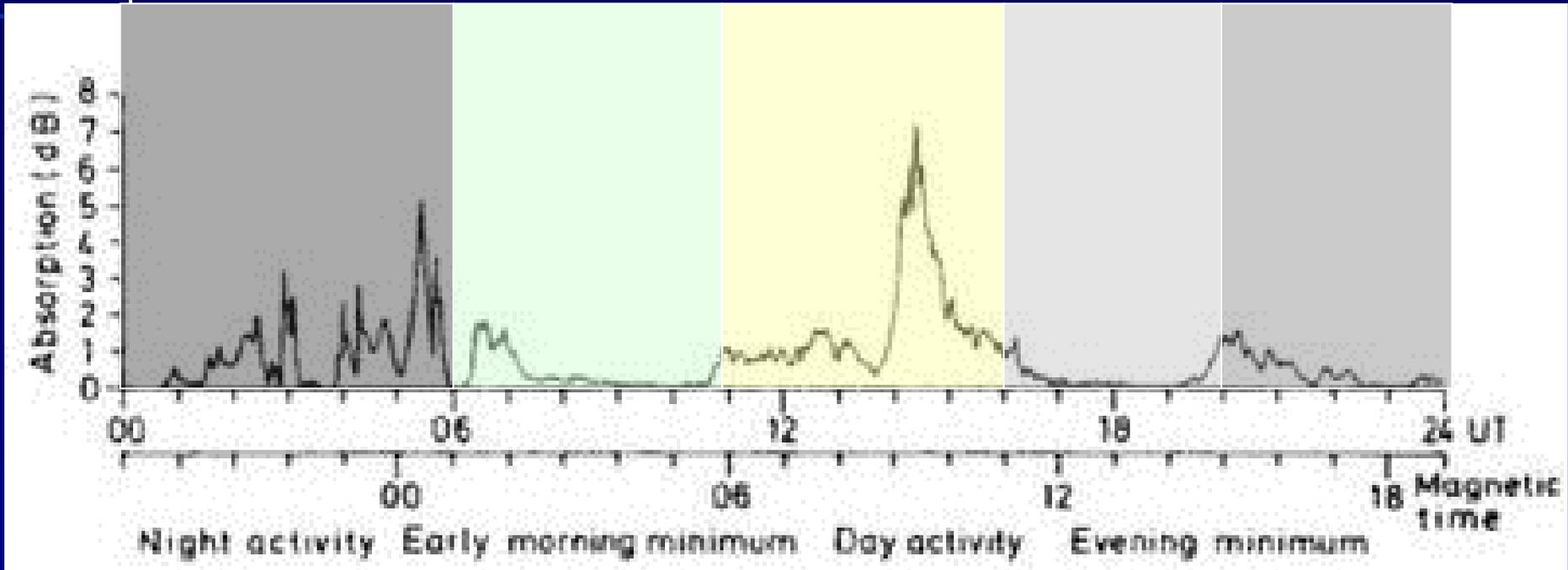


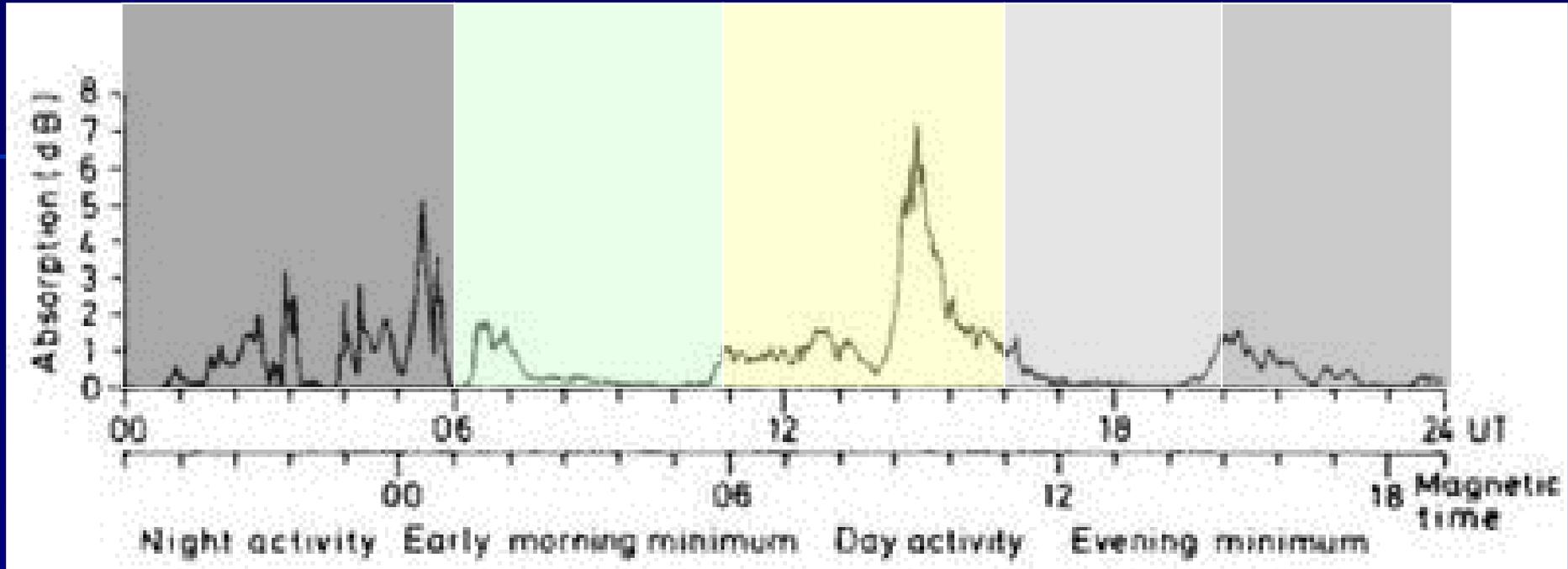
Chart of Attenuation

Carrier radio wave frequency should be change with time ! Because absorption rate of radio waves changing with time in the following diagram.



Ionosphere absorption (or ISAB) is only a factor in the period of the day where radio signals travel through the portion of the ionosphere facing the **Sun**. The **solar wind** and **radiation** cause the ionosphere to become charged with electrons in the first place.

Chart of Attenuation



At night, the atmosphere become drained (fall) of its charge, and radio signals can go much further with less loss of signal. In particular, low wavelength signals that would be attenuated to nothing during the day will be received much further away at night.

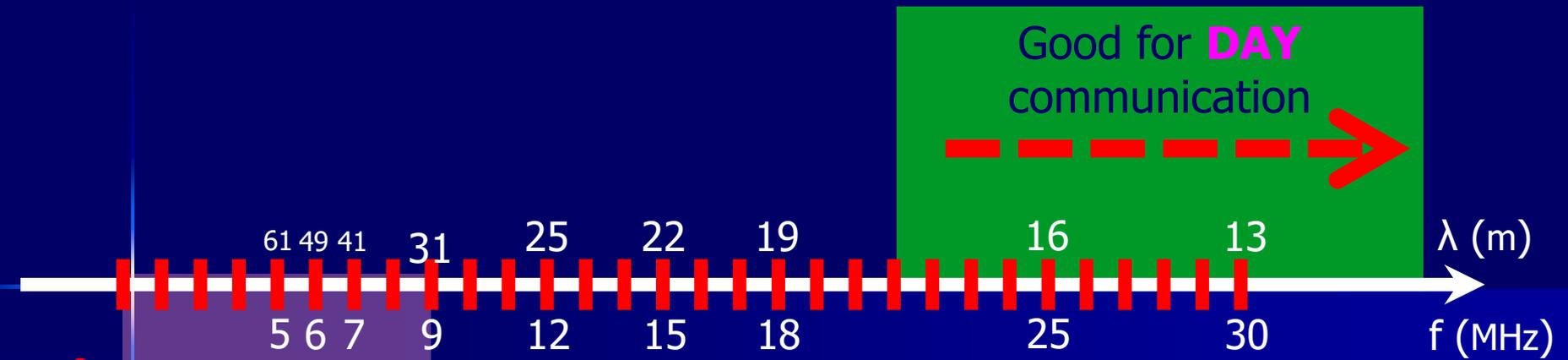
The specific amount of attenuation can be derived as a function of the inverse square law. The lower the frequency, the greater the attenuation.



In radio wave communication, the **Radio Broadcasting Center** should change their carrier frequencies in the day and the night like the above diagram. Also these frequencies depend on the number of sunspots and activity of the Sun.

Eg: **BBC Sandeshaya**

(They used the transmission frequency **250 kHz** for the communication during the **DAY** and transmission frequency **100 kHz** for the **NIGHT**)



BBC Sandeshaya (NOW)

Daily **BBC Sandeshaya** program is broadcasting on SW 19 (15,690 kHz), SW 22 (13630 kHz) and SW 41 (7435 kHz) at 9:15 pm (night) on Sri Lankan time !

<http://www.bbc.com/sinhala>

Generally India using **SW communication** to broadcast their radio channels for all over the country !

Radio Wave Communication

Radio waves

Radio Communication

Reflection of Radio Waves

Absorption of Radio Waves

Complex Refractive Index

Reflection Heights

Deviating Region Absorption, Non-

Deviating Region Absorption

Ordinary/Extra Ordinary Waves

Ionosphere – Sounding Techniques

Pulse Reflection Methods

Complex Refractive Index

If the radio wave going through the ionosphere, **refractive index n is always a complex number**. (As we discussed in the ionosphere chapter)

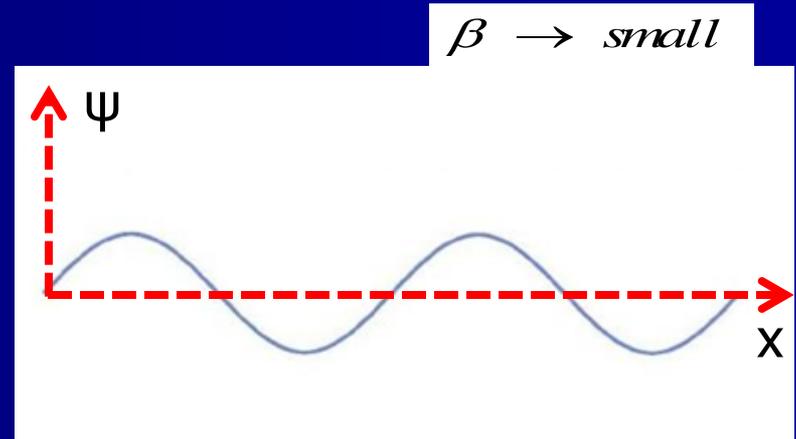
That means $n = \alpha + i \beta$ and n is complex. Therefore, we can use $n \rightarrow n^*$. i.e.;

$$n^* = \alpha + i \beta$$

Where α and β are positive constants. And $i = \sqrt{-1}$

\therefore Equation of the motion of the radio wave can be express as the following form;

$$\psi = \psi_0 e^{i n x}$$



Complex Refractive Index

$$\psi = \psi_o e^{i n x}$$

Now we can substitute, $\mathbf{n = n^* = \alpha + i \beta}$ to the above equation.

$$\therefore \psi = \psi_o e^{i(\alpha+i\beta)x}$$



$$\psi = \psi_o e^{i\alpha x} e^{-\beta x}$$



$$\psi = (\psi_o e^{-\beta x}) e^{i\alpha x}$$



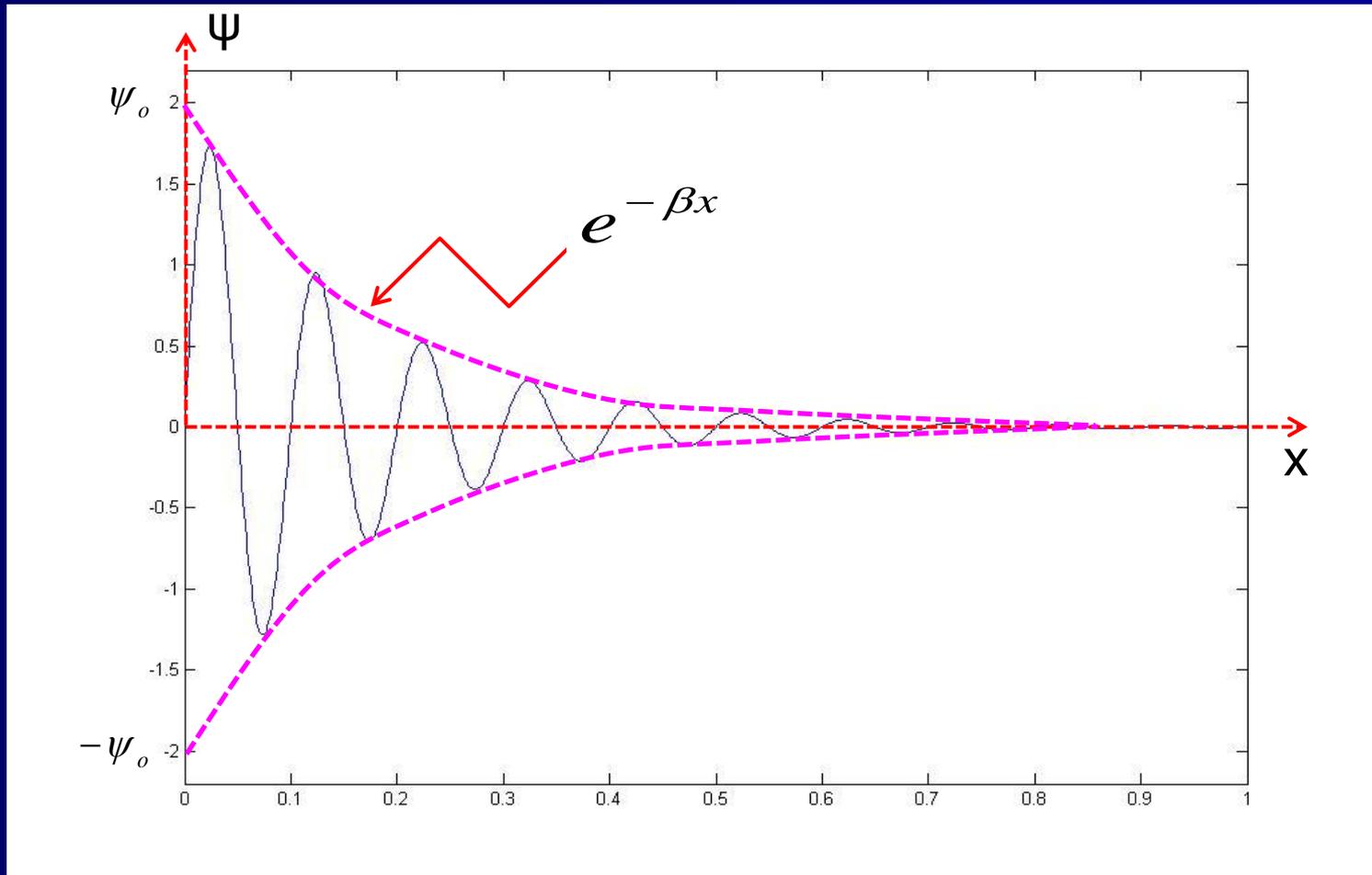
$$\psi = \psi_o^1 e^{i\alpha x} \quad \text{and} \quad \psi_o^1 = \psi_o e^{-\beta x}$$

Where $\psi_o^1 = \psi_o e^{-\beta x}$ is the new amplitude of the Radio Wave.

Complex Refractive Index

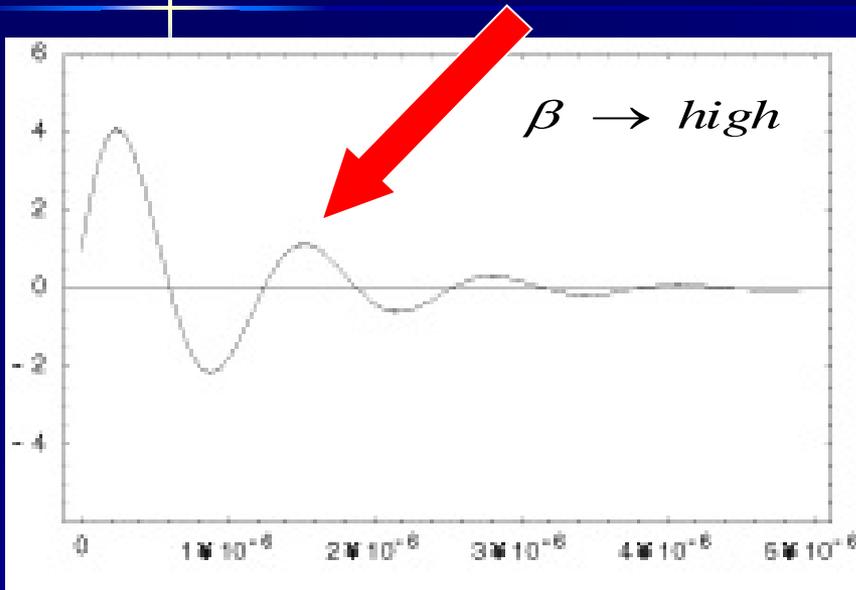
$$\psi_o^1 = \psi_o e^{-\beta x}$$

means that the new amplitude is attenuated from the ionosphere.

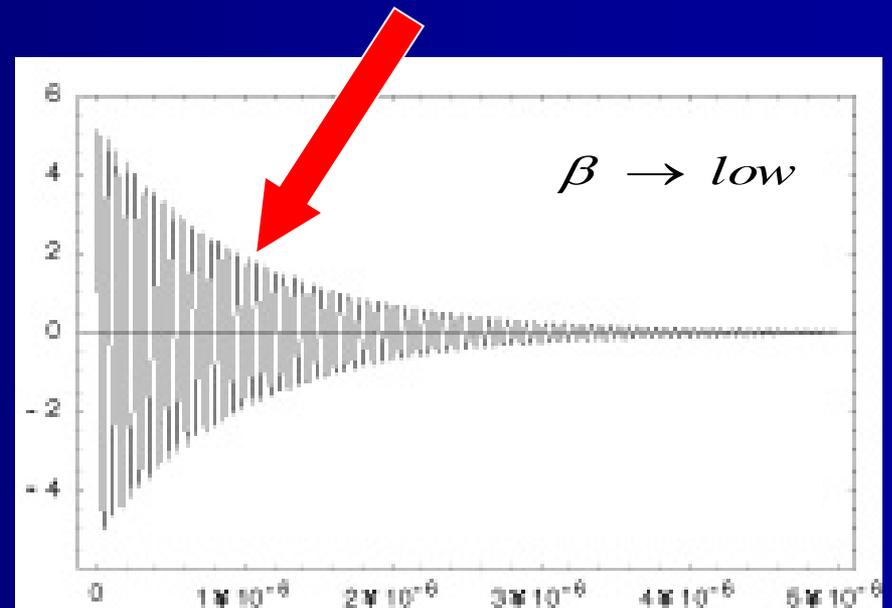


Complex Refractive Index

$$n = \alpha + i\beta$$



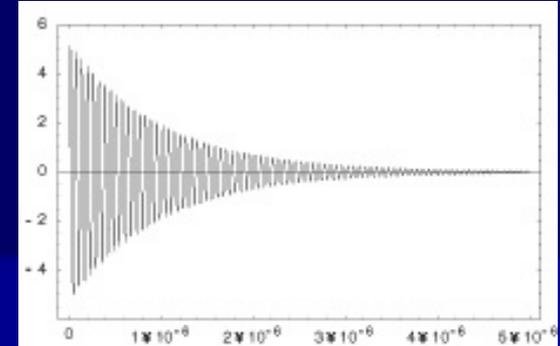
Attenuation



Complex Refractive Index

The wave is attenuated. That means the wave is decaying. Because, the wave is absorbed from the ionosphere.

$$Ab \propto \frac{1}{c + f}$$



- If the **frequency of the radio wave is very high**, the **absorption from the ionosphere is very low**. That means the **attenuation of the radio wave is small**. That means **Exp[-βx]** is **approximately equal to one**.

$$e^{-\beta x} \approx 1$$



$$\beta \rightarrow 0$$

That means complex part of the refractive index ($n^* = \alpha + i\beta$) is negligible. It is going through the ionosphere.

Complex Refractive Index

- If the **frequency of the radio wave is very small**, the absorption from the ionosphere is very high. That means the **attenuation of the radio wave is high**. That means **Exp[-βx]** is going to zero.

$$e^{-\beta x} \approx 0$$

 **β is very large !**

That means complex part of the refractive index ($n^* = \alpha + i\beta$) is very large. It is reflected back from the ionosphere to the Earth surface.

Plasma frequency of the ionosphere: $f_p = 8.97 \times N^{1/2}$

If the electron density of the ionosphere is $10^{12} \text{ e}^{\text{n}}/\text{cm}^3$,

$$f_p = 8.97 \times (10^{12})^{1/2}$$

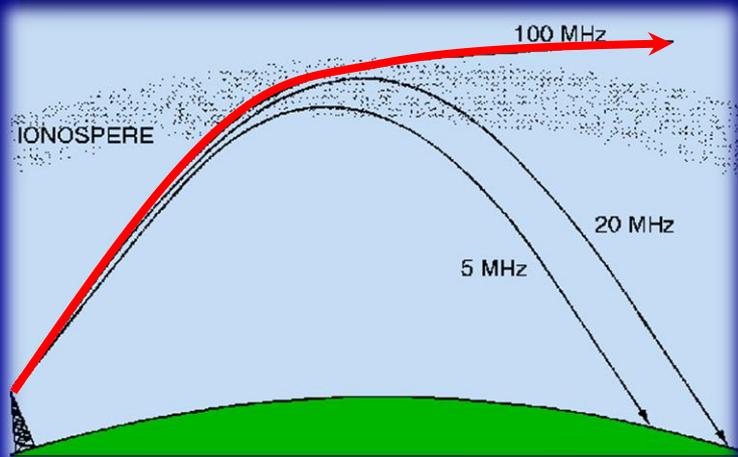
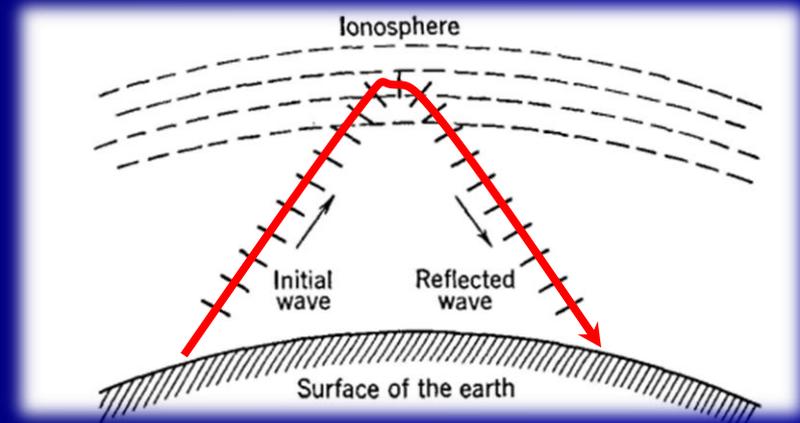


$$f_p \approx 9 \text{ MHz}$$

Complex Refractive Index

- That means, if we send a **Radio Wave** to the ionosphere with frequency **9 MHz** (or lower value), it is **reflected from the region of the ionosphere** whose electron density is $10^{12} \text{ e}^n/\text{cm}^3$.

That means **VLF** (3-30 kHz), **LF** (30-300 kHz), **MF** (300-3000 kHz), and **HF** (some short waves in the region 3 MHz - **9 MHz**) are **reflected from the ionosphere** !



But, if we send **UHF** (>30 MHz) and **VHF** (>300 MHz) signal to the ionosphere, both signals are **going through ionosphere** to the space without reflection back!

Complex Refractive Index

Complex index of refraction is $\mathbf{n = n^* = \alpha + i\beta}$, in looking in detail at the **polarization of dielectrics**, we switched from a simple di-electric constant ϵ_r to a **Complex di-electric function** $\epsilon_r(\omega) = \epsilon_1 + i\epsilon_2$.

i.e.; $\epsilon_r(\omega) = \epsilon_1 + i\epsilon_2$

Also we know; $n^* = \alpha + i\beta$

With the **di-electric constant** ϵ_r and **refraction index constant** n^* , we had the basic relation;

$$n^2 = \epsilon_r \quad \rightarrow \quad n^{*2} = \epsilon_r \quad \rightarrow \quad (\alpha + i\beta)^2 = \epsilon_1 + i\epsilon_2$$

That means, $\alpha = \alpha(\omega)$ and, $\beta = \beta(\omega)$

$$\rightarrow (\alpha + i\beta)^2 = \epsilon_1 + i\epsilon_2 \quad \rightarrow \quad \alpha^2 - \beta^2 + 2(\alpha\beta)i = \epsilon_1 + i\epsilon_2$$

Complex Refractive Index

$$\alpha^2 - \beta^2 + 2(\alpha\beta)i = \varepsilon_1 + i\varepsilon_2$$

→ $\alpha^2 - \beta^2 = \varepsilon_1$ and, $2\alpha\beta = \varepsilon_2$

.....

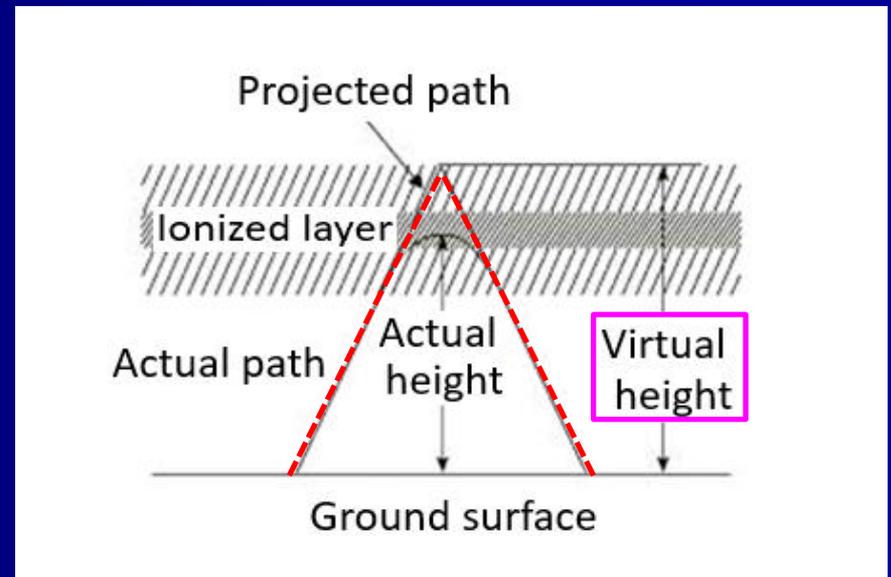
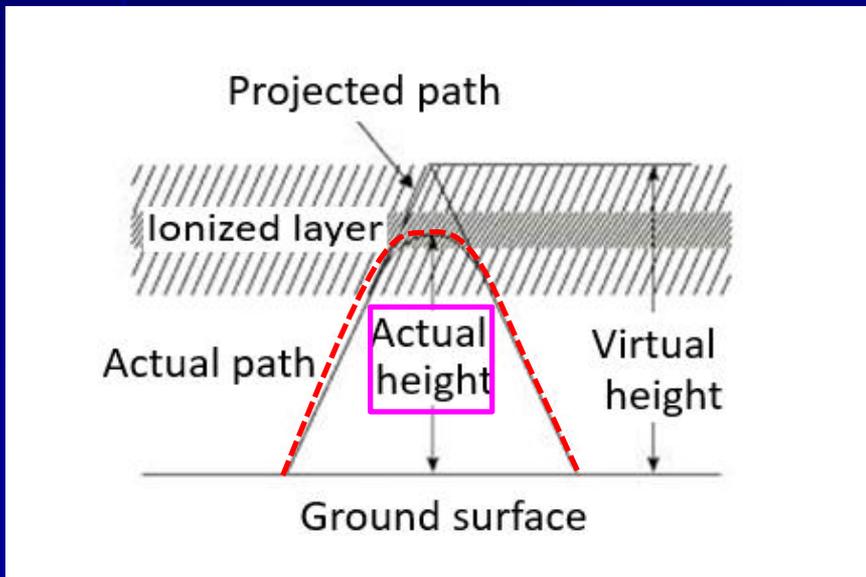
→ $\alpha^2 = \frac{1}{2} \left[(\varepsilon_1^2 + \varepsilon_2^2)^{1/2} + \varepsilon_1 \right]$ and, $\beta^2 = \frac{1}{2} \left[(\varepsilon_1^2 + \varepsilon_2^2)^{1/2} - \varepsilon_1 \right]$

Using above equations we can find the values of α and β when ε_1 and ε_2 are already known.

- If $\beta \rightarrow 0$ then $\varepsilon_2 \rightarrow 0$
that means $\varepsilon_r(\omega) = \varepsilon_1$. (Real)
- If $\beta \rightarrow \infty$ (very large) then $\varepsilon_2 \rightarrow \infty$ (very large)
that means $\varepsilon_r(\omega) = i\varepsilon_2$. (complex)

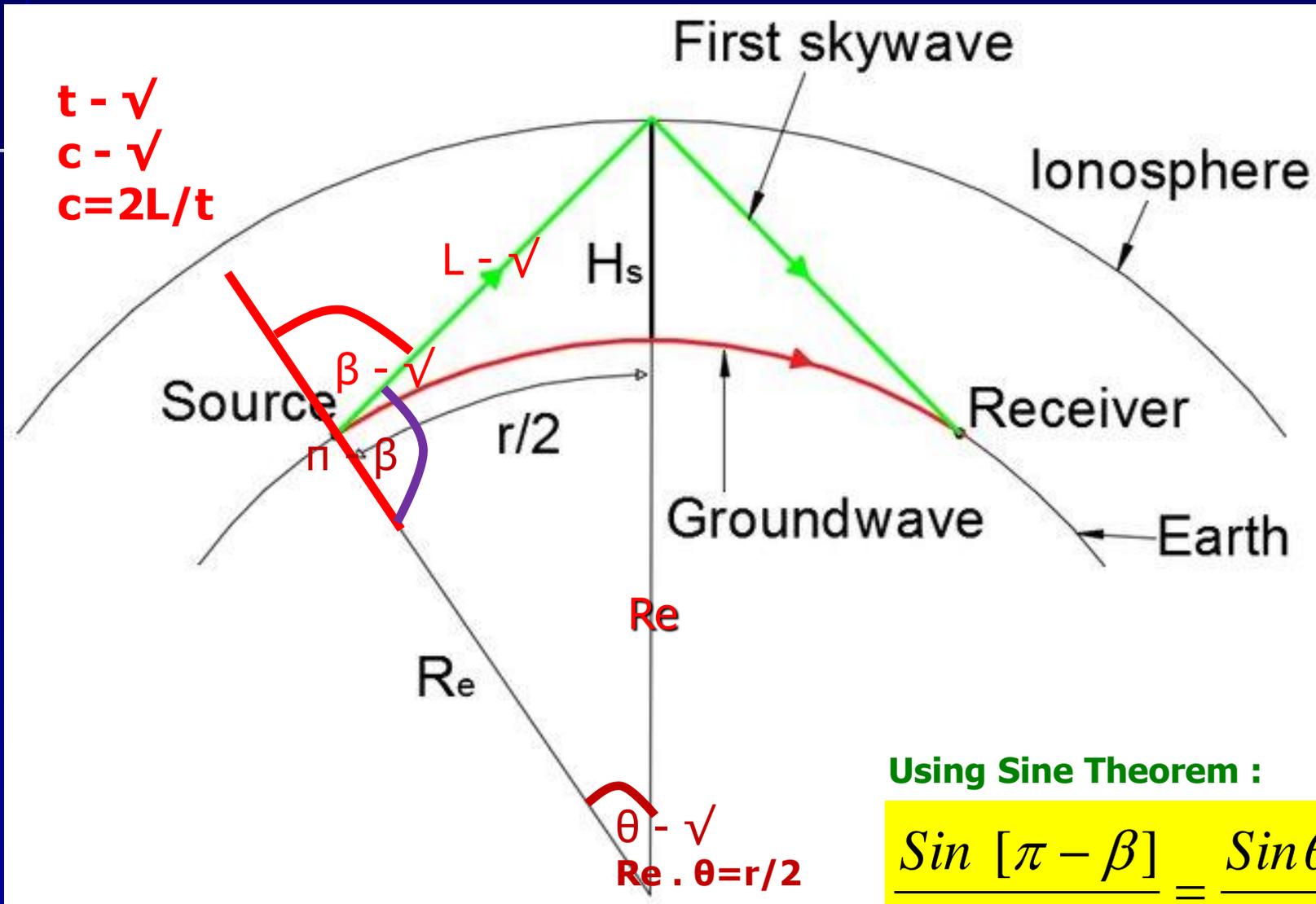
Reflection Heights

A radio wave transmitted towards the ionosphere is subject to be reflected. The process of this reflection shows an exciting feature which is, that the position it reflects from the ionosphere is varying with corresponding to the frequency of a particular radio wave. This particular position of reflection is expressed in terms of heights and it is known as **reflection height** for the particular radio wave.



Reflection Heights

Incident angle - ν



$$t - \nu$$

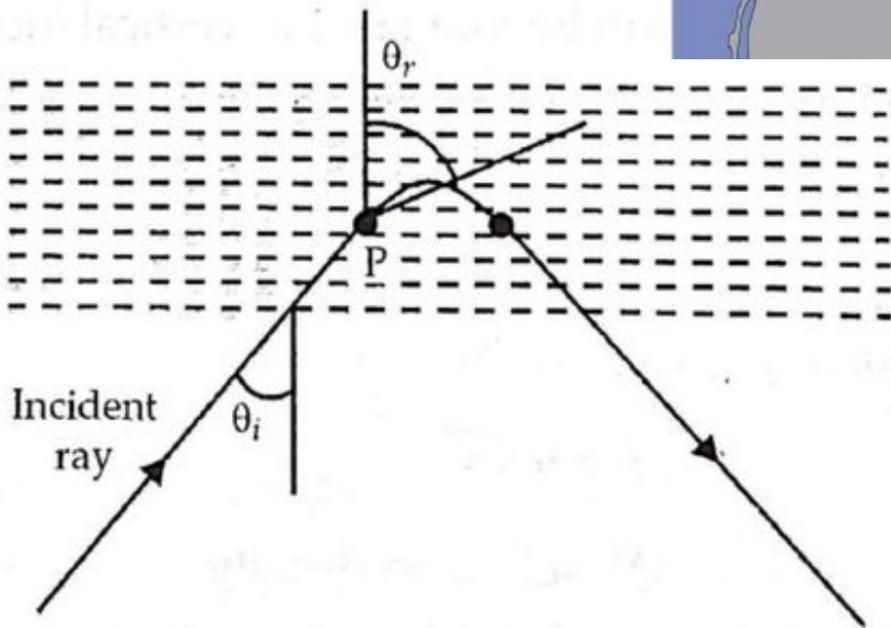
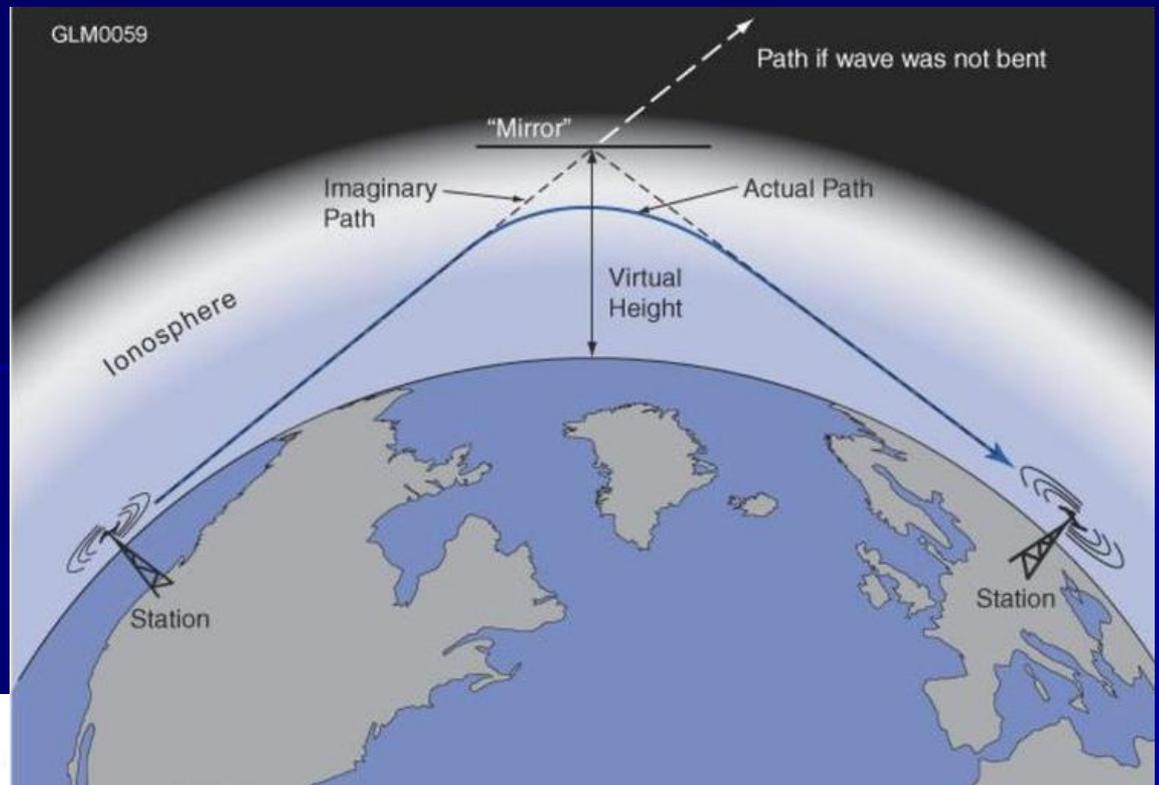
$$c - \nu$$

$$c = 2L/t$$

Using Sine Theorem :

$$\frac{\sin [\pi - \beta]}{R_e + H_s} = \frac{\sin \theta}{L}$$

Reflection Heights



Snell's Law of refraction in the ionosphere.

$$n = \frac{\sin \theta_i}{\sin \theta_r}$$

n - is refractive index

θ_i - is angle of incident

θ_r - is angle of refraction

$$H(r, \theta) = \frac{\mu_0}{4\pi} \cdot \frac{M}{r^3} \cdot (1 + 3 \cos^2 \theta)^{\frac{1}{2}}$$

Reflection Heights

$$f_C = \sqrt{81 N_{max}} = 9 \sqrt{N_{max}}$$

$$f_C = \sqrt{81 N_{max}}$$

$$f_C^2 = 81 N_{max}$$

$$\frac{81 N_{max}}{f_C^2} = 1$$

$$1 - \frac{81 N_{max}}{f_C^2} = 0$$

$$n = \sqrt{\epsilon_r} = \sqrt{1 - \frac{81 N}{f^2}}$$

We know $n = \sqrt{1 - \frac{81 N}{f^2}}$

$$n = \frac{\sin \theta_i}{\sin \theta_r} \rightarrow \frac{\sin \theta_i}{\sin \theta_r} = \sqrt{1 - \frac{81 N}{f^2}}$$

Angle of incident $\theta_i = 0$

when $\theta_i = 0$, $N = N_{max}$ and $f = f_C$

Therefore, $0 = \sqrt{1 - \frac{81 N_{max}}{f_C^2}}$

Reflection Heights

Now, the condition for the wave to be reflected back is given by

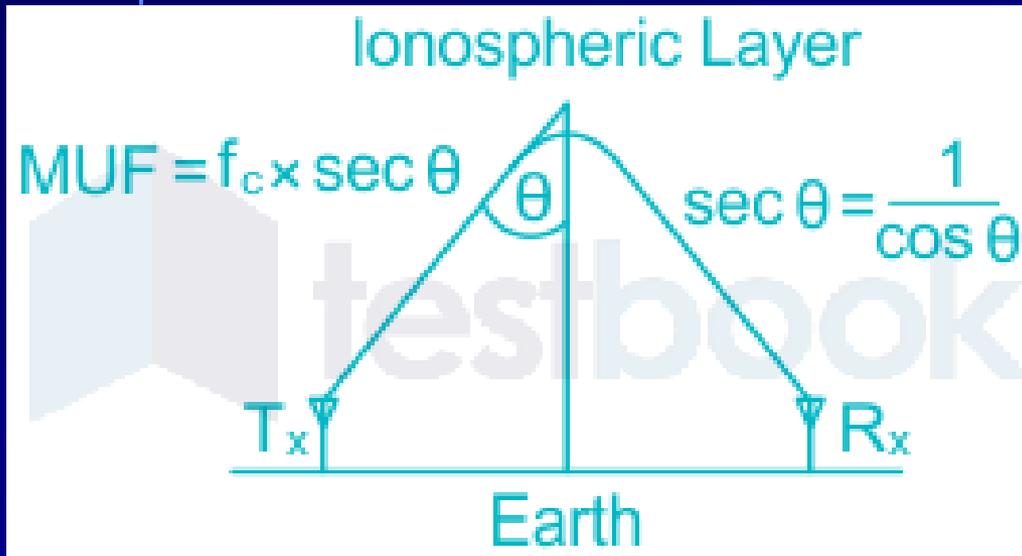
$$\sin \theta_i > \sqrt{1 - \frac{81 N_{max}}{f^2}}$$

$$\text{But } f_c = \sqrt{81 N_{max}}$$

$$\therefore \sin \theta_i > \sqrt{1 - \left(\frac{f_c}{f}\right)^2}$$

Maximum usable frequency (MUF) :

MUF (Maximum Usable Frequency) is the maximum frequency which can be reflected for given distance of transmission. MUF is usually 3 to 4 times of critical frequency.



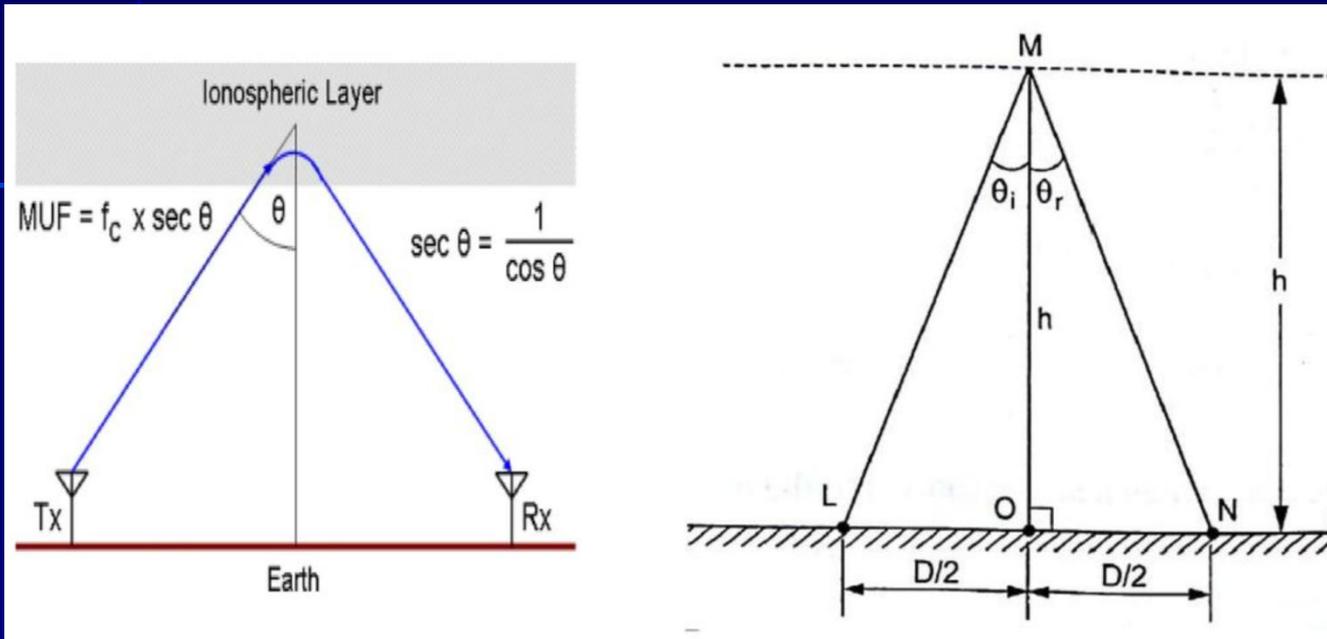
$$MUF = \frac{\text{critical frequency}}{\cos \theta}$$

Where,

MUF = Maximum Usable Frequency

θ = Angle of Incidence

Reflection Heights



Considering ΔLOM in Fig.

$$\cos \theta_i = \frac{h}{LM}$$

$$= \frac{h}{\sqrt{\left(\frac{D}{2}\right)^2 + h^2}}$$



$$\cos \theta_i = \frac{2h}{\sqrt{D^2 + 4h^2}} \longrightarrow 1$$

Reflection Heights

We know

$$\therefore n = \sin \theta_i = \sqrt{1 - \frac{81 N_{max}}{f_{MUF}^2}}$$

$$\sin^2 \theta_i = 1 - \frac{81 N_{max}}{f_{MUF}^2} \Rightarrow 1 - \sin^2 \theta_i = \frac{81 N_{max}}{f_{MUF}^2}$$

$$\cos^2 \theta_i = \frac{81 N_{max}}{f_{MUF}^2}$$

We know $f_c = \sqrt{81 N_{max}}$

$$\therefore \cos^2 \theta_i = \frac{f_c^2}{f_{MUF}^2}$$

Comparing Equ. 1 & 2.

$$\frac{2h}{\sqrt{D^2 + 4h^2}} = \frac{f_c}{f_{MUF}}$$

$$f_{MUF} = f_c \left[\frac{\sqrt{D^2 + 4h^2}}{2h} \right]$$

$$f_{MUF} = f_c \sqrt{\left(\frac{D}{2h}\right)^2 + 1}$$

$$f_{MUF} > f_c$$

$$\therefore \boxed{\cos \theta_i = \frac{f_c}{f_{MUF}}} \longrightarrow 2$$

Reflection Heights

- The Skip distance, d_s is

$$d_s = \frac{2h}{\tan \theta_c}$$

h – height of the layer

θ_c – Critical angle

A **skip distance** is the distance a radio wave travels, usually including a hop in the ionosphere. A skip distance is a distance on the Earth's surface between the two points where radio waves from a transmitter, refracted downwards by different layers of the ionosphere, fall.

We know

$$f_{\text{MUF}} = f_c \sqrt{\left(\frac{D}{2h}\right)^2 + 1}$$

$$f_{\text{MUF}} = f_c \sqrt{1 + \left(\frac{D_{\text{skip}}}{2h}\right)^2}$$

$$\left(\frac{f_{\text{MUF}}}{f_c}\right)^2 = 1 + \left(\frac{D_{\text{skip}}}{2h}\right)^2$$

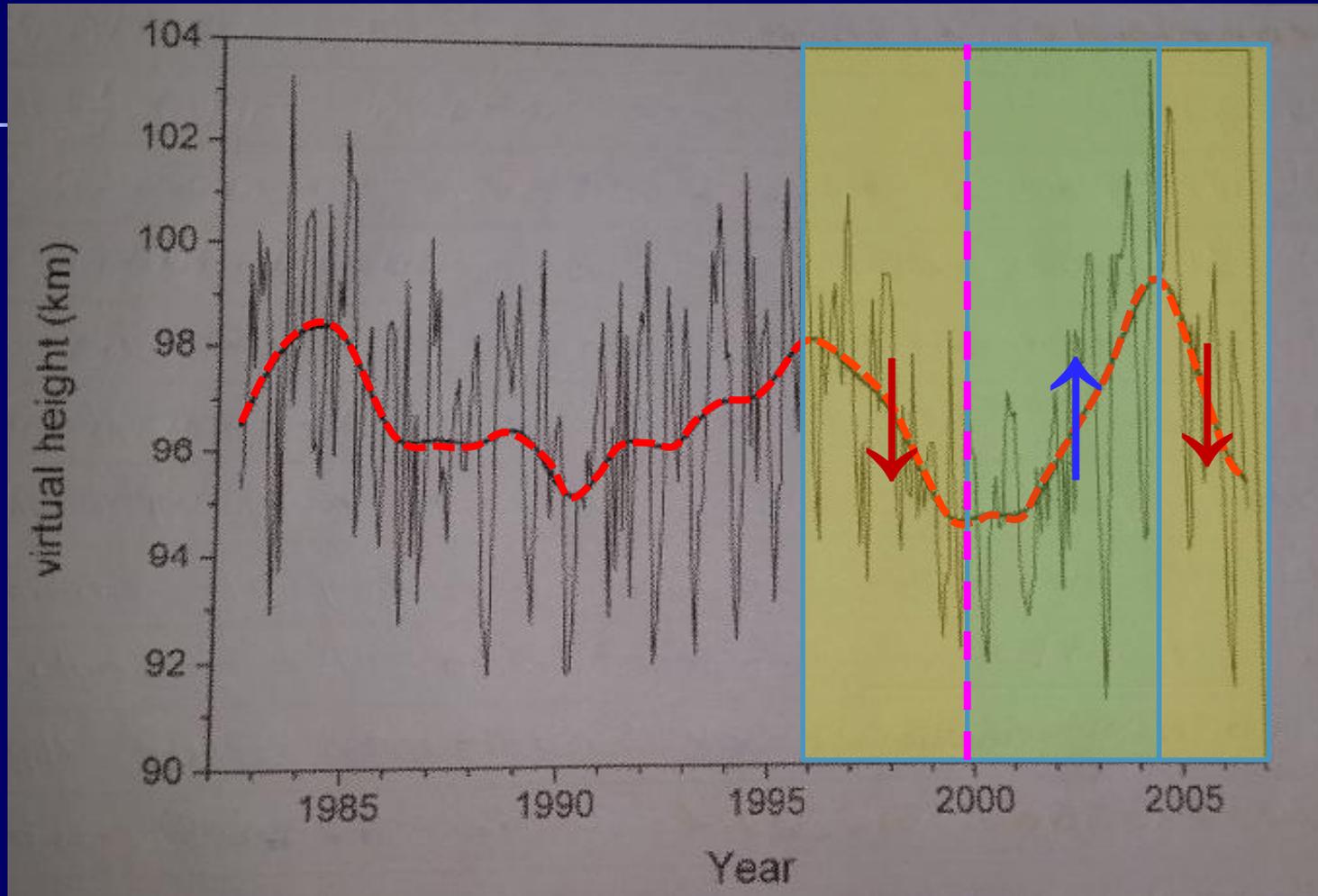
$$D_{\text{skip}} = 2h \sqrt{\left(\frac{f_{\text{MUF}}}{f_c}\right)^2 - 1}$$

Reflection Heights

Trends in low frequency radio wave reflection heights and lower E-region drifts over Collm (a city), Germany.

Lower E-region low frequency measurements using commercial radio broadcasting stations have been used to indirectly measure mesospheric temperature variability via change of the radio wave reflection height. Recent analyses have shown that the temperature decrease is obviously not linear, but ceases (stops) during the last decade. Collm (a city in Germany) measurements of reflection heights on **177 kHz**, referring to a reflection point at **(52.1°N,13.2°E)** are shown in the following figure. These virtual heights are about 5 km too high compared with the real ones. **Until about year 2000, the reflection height decreases**, but after that time the lower E-region heights remain at a similar level.

Reflection Heights



Collm measurements of reflection heights on **177 kHz**, referring to a reflection point at **(52.1°N,13.2°E)** are shown in the above figure.

Reflection Heights

Low frequency measurements are also used to analyze E-region drifts can be interpreted as neutral winds at about 90 km altitude. Analyses of long-term trends have **revealed** possible long-term trends in some parameters, generally through to be in **connection with middle atmosphere changes**. Semi-diurnal (half daily) tidal amplitude have decreased, possibly in connection with **stratospheric ozone decrease**. **Meridional** (southern) **winds have decreased** due to **stratospheric cooling**, decrease of mesospheric wind jets and a resulting weaker gravity wave filtering which leads to **reduced Brewer-Dobson circulation**.

Generally, there is an increase of the zonal prevailing wind, a decrease of Meridional wind, which means a positive trends in summer (negative northerly winds) and a negative trend in winter (southerly weak wind). The trends are not linear. The strongest trends visible in the early 1990's while before and after that time trends are weak, and during recent years, even seem to reverse. The Meridional winds have decreased. In particular the summer trend weakens. The semi-diurnal tidal amplitudes show a complicated behavior, there are large values in the early 1990's, possibly connected with the **11 year solar cycle**.

Reflection Heights

To **conclude**, the long-term wind and reflection height time series show trends during the last 3 decades. But there seem to be changes of these trends after the 1990's. This is clearly visible in the prevailing winds and also in temperatures. It is generally thought that middle atmosphere climate change is responsible for at least part of the lower E-region changes. This means, that in addition to the continuing CO₂ cooling other influences act on the upper middle atmosphere. One candidate is Ozone, which has decreased strongly since 1980's, but had remained at a constant level or even recovered after the mid 1990's.

Radio Wave Communication

Radio waves

Radio Communication

Reflection of Radio Waves

Absorption of Radio Waves

Complex Refractive Index

Reflection Heights

Deviating Region Absorption, Non-Deviating Region Absorption

Ordinary/Extra Ordinary Waves

Ionosphere – Sounding Techniques

Pulse Reflection Methods

Deviation Region Absorption & Non-Deviation Region Absorption

There are many different type of phenomena that attenuate or reduce the signal strength of propagating radio waves. Those phenomena types extract energy from the radio wave by converting its energy to **heat** and **electromagnetic noise** and are associated with the term absorption.

When referring to the actual propagation of a radio wave through the ionosphere, we have **principally two types of ionospheric absorption.**

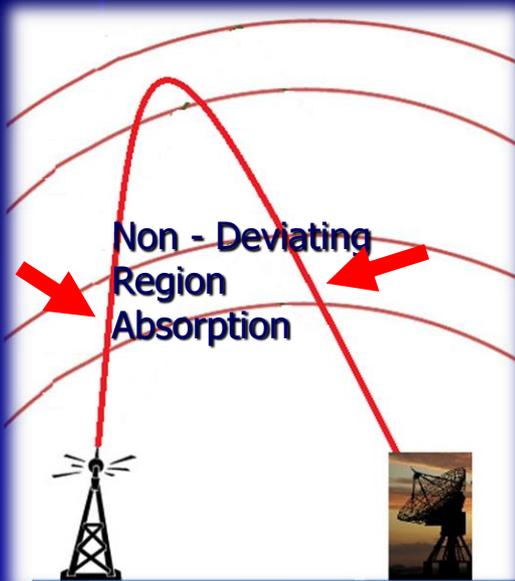
01. Deviating Absorption

02. Non - Deviating Absorption

During the day the ionospheric electron density is much increased but recombine also at a quick rate with the ever-present high densities of ions and neutral particles, especially in the lower ionosphere where the gasses are denser than at the higher more verified regions. **This produces stronger levels of ionospheric non-deviating absorption.**

Deviation Region Absorption & Non-Deviation Region Absorption

In contrast, during the night, when the electron densities are substantially lower, the non-deviating absorption decreases. The following equation summarizes empirically the **non-deviating absorption**:



*Non-Dev Region
Absorption Coefficient*

$$\kappa = 1.15 \times 10^{-3} \frac{N \nu}{f}$$

*Ele; Density in the
Lower Ionosphere*

Freq; of the Radio wave

*Collision
Frequency of ele;
with neutral
particles*

Deviation Region Absorption & Non-Deviation Region Absorption

$$\kappa = 1.15 \times 10^{-3} \frac{N \nu}{f}$$

From the above equation, we can see that; if the electron density **N** or the collision frequency **ν** is increased, while the radio frequency **f** stays constant, then the absorption level **κ** will increase,

If **N** ↑ or **ν** ↑ → **κ** ↑ (when **f** constant)

However, if the radio frequency (**f**) is increased then the effect of non-deviating absorption (**κ**) decreases.

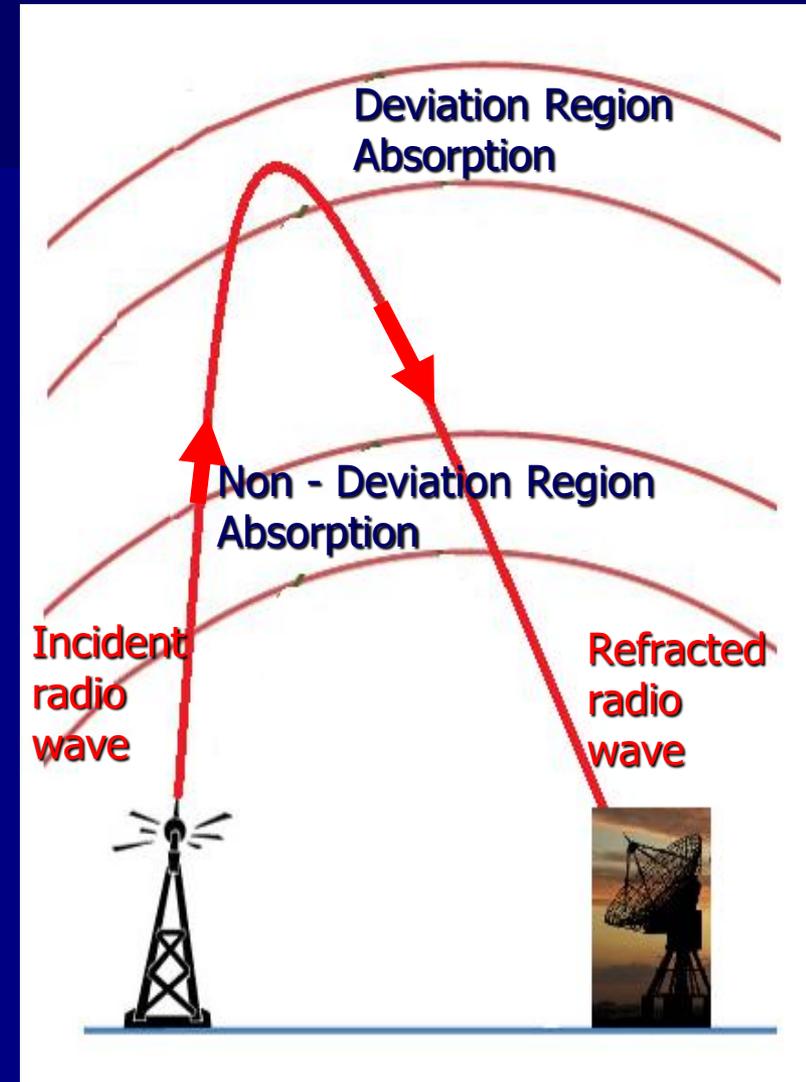
If **f** ↑ → **κ** ↓

Lower frequencies have a greater or larger wave front than higher frequencies (even when the amplitude is equal) !

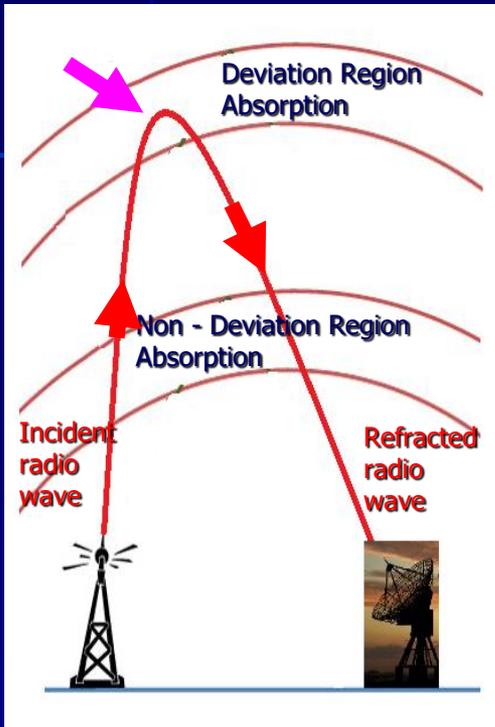
Deviation Region Absorption & Non-Deviation Region Absorption

A greater wave front will collide much more with ions and neutral particles and therefore be more attenuated. If we send a wave to the ionosphere, the wave is spending the additional time interval in the ionosphere due to more collisions with ions and neutral particles.

That extra spending time interval is large, that means, the wave meet Deviation Region or not. If we analyzed the time interval of the incident and reflected waves, we can conclude the wave meet a peak (deviating region) or not.



Deviation Region Absorption & Non-Deviation Region Absorption



Sometimes **electron density** of the ionosphere is having a **very large values** when **huge solar activity**, **VHF** (30 MHz – 300 MHz) and **UHF** (300 MHz – 3 GHz) **also may be reflected** from the ionosphere.

Deviative Absorption is a type of absorption that occurs when the signal spends a **longer period** of time within the absorbing medium and meanwhile is refracting. The derivative absorption coefficient depends heavily upon the **refractive index of the signal** and frequency of the electron collisions with neutral particles.

*Deviation Region
Absorption Coefficient*

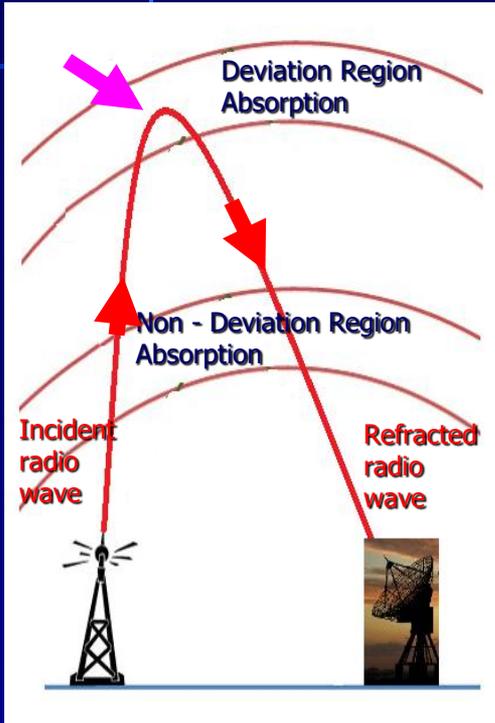
$$\kappa^1 = \kappa^1(n, f)$$

*Collision Frequency
of ele; with neutral
particles*

*Refractive Index
of the Signal*

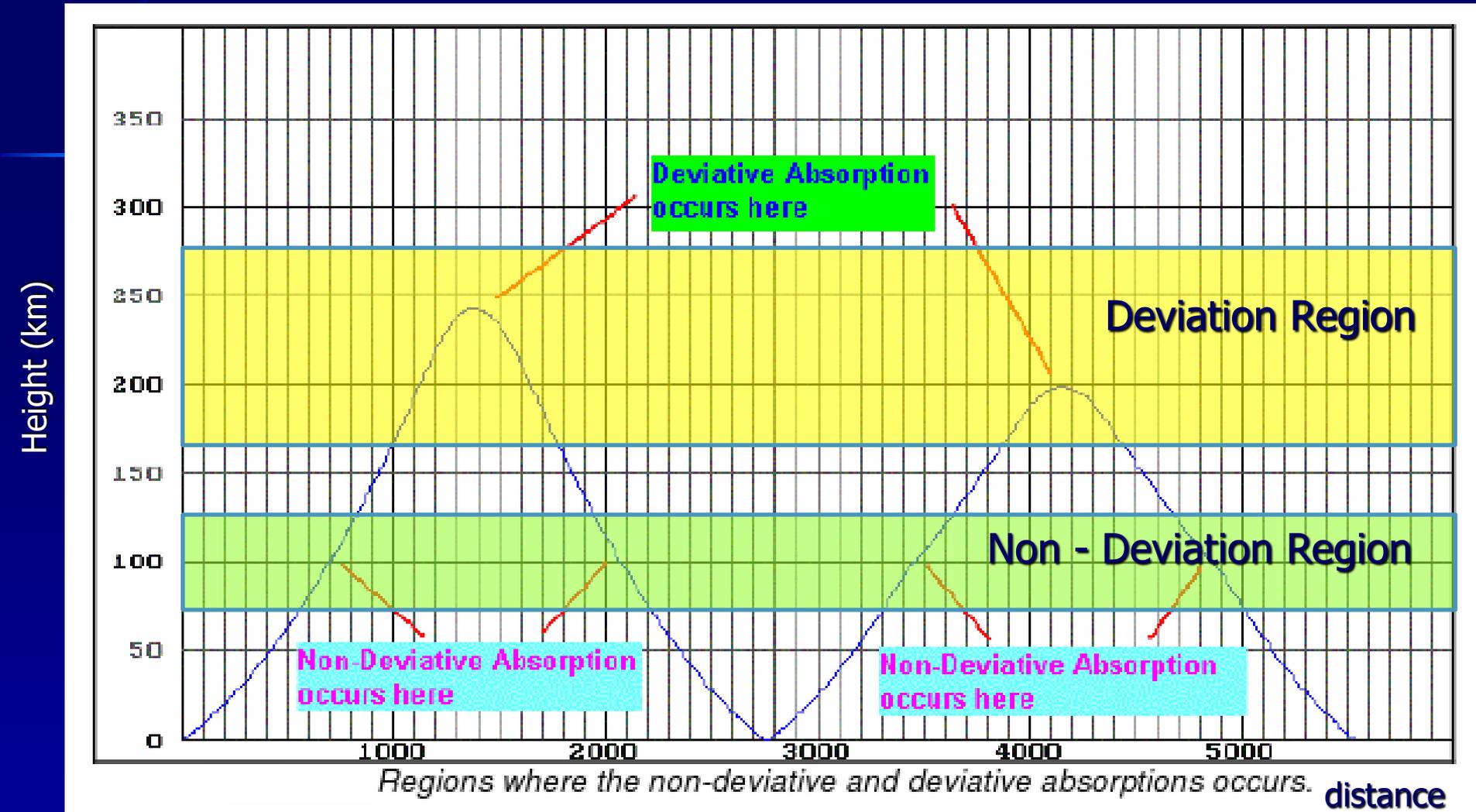
Deviation Region Absorption & Non-Deviation Region Absorption

$$\kappa^1 = \kappa^1(n, f)$$



Each time a signal is refracted in the ionosphere, deviating absorption occurs. Refraction in the lower ionospheric regions, where the electron collision frequency is rather high, experience greater absorption and can be rather high, particularly the frequencies in the lower **HF** spectrum. If the **Sun** also **illuminates** (very active) the communication circuit, then the effects of non-deviating absorption will take an additional toll on (badly damaged) the signal strength.

Deviation Region Absorption & Non-Deviation Region Absorption



Deviating absorption is usually not as strong as Non-Deviating absorption !

Radio Wave Communication

Radio waves

Radio Communication

Reflection of Radio Waves

Absorption of Radio Waves

Complex Refractive Index

Reflection Heights

Deviating Region Absorption, Non-

Deviating Region Absorption

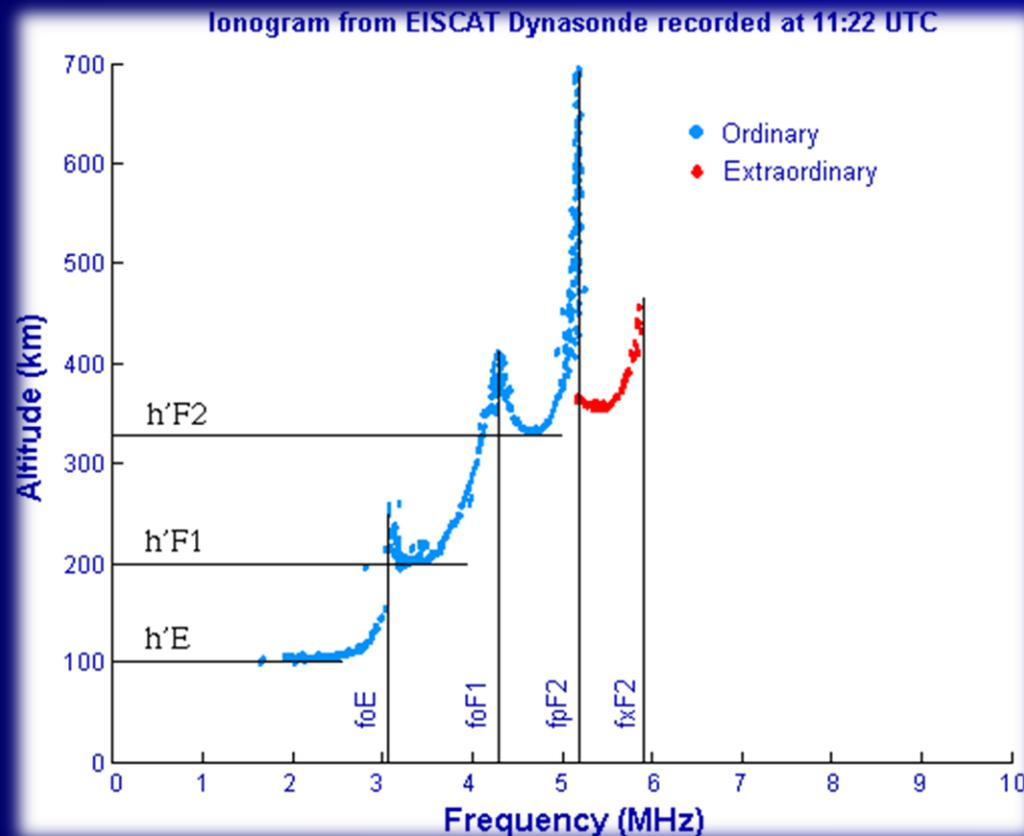
Ordinary/Extra Ordinary Waves

Ionosphere – Sounding Techniques

Pulse Reflection Methods

Ordinary & Extraordinary WAVES

As soon as transmitted radio signals begins to penetrate the ionosphere, **it splits up into two characteristic components due to the Earth's Magnetic Field**. These two components are known as the **ordinary** and **extraordinary** waves.



Ordinary & Extraordinary WAVES

For the **ordinary wave**, the E-field accelerates electrons **parallel** to the **magnetic field**. This means that the electric field has no influence, because a magnetic field only imposes a force on charged particles moving perpendicular to the Field !

For the **extraordinary wave**, the E-field of the incident radiation accelerates the free electrons **normal** to the magnetic field. This means that it exerts a force on the electrons and therefore modifies the motion. This causes the **refractive index of the extraordinary wave** to be **different from the ordinary wave**.

The **different refractive indices** of the two component waves, meaning **different velocities**, also cause a **progressive phase shift** between the two components. If the **phase shift becomes 90°** , then the initial **100% circular polarized wave**. For **smaller different phase shifts** will the wave become **elliptically polarized**.

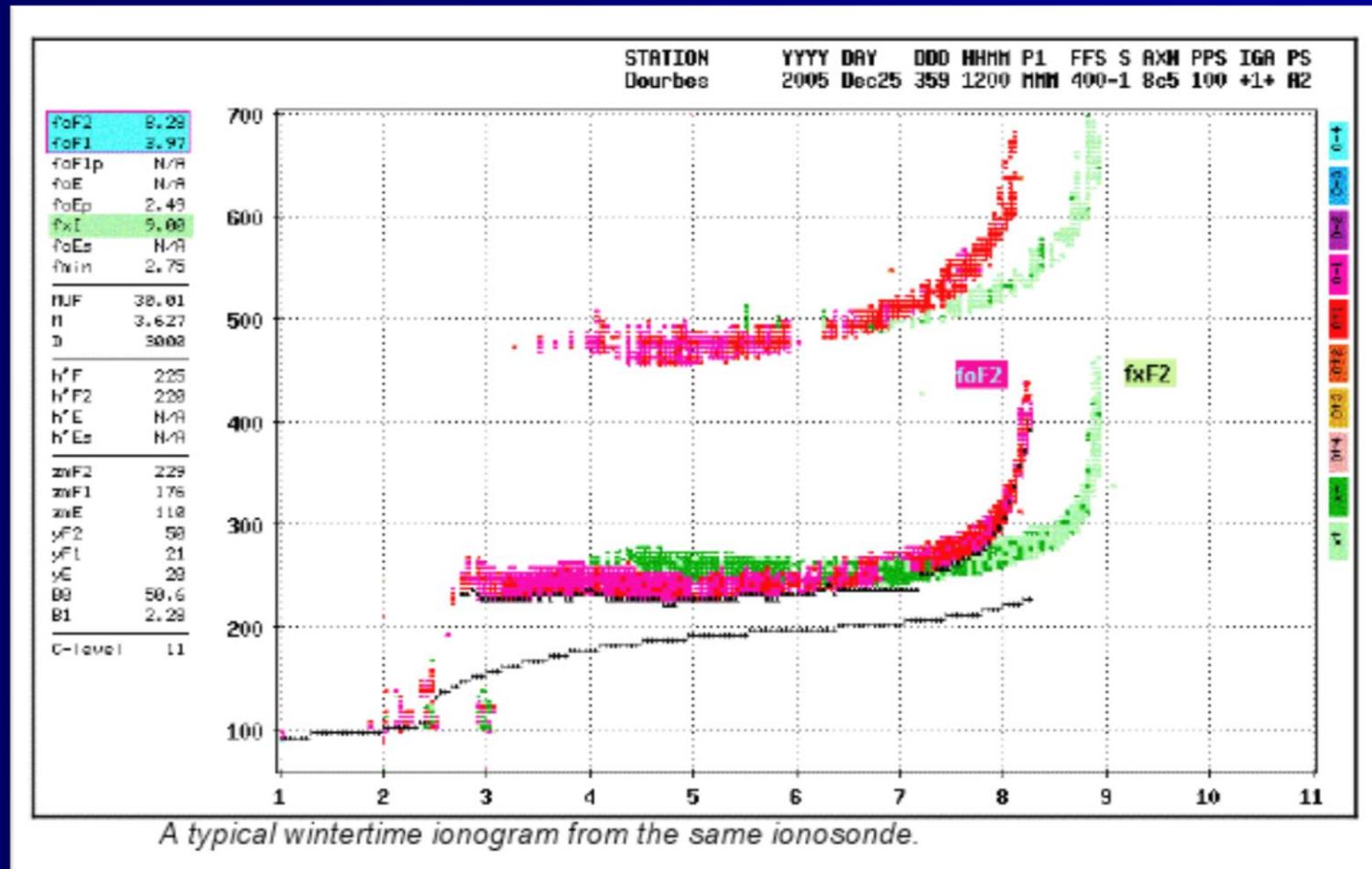
Ordinary & Extraordinary WAVES

Each of these waves, the ordinary and the extraordinary identified as the o-wave or the x-wave, travels a completely independent path through the ionosphere. The single radio signal travels as simultaneously transmitted but independent transmissions. Each propagating signal component contributes to a different mixture of power level, which totals together the power the signal before it enters the ionosphere. **The individual power level depends upon some complicated relationships between the o-wave and x-wave at the base of the ionosphere,** but the x-wave is the weaker of the two.

- At **higher frequencies** the o-wave and x-wave often follow very simple path.
- At **lower frequencies** the o-wave and x-wave will diverge more considerably.

Ordinary & Extraordinary WAVES

The existence of the ordinary waves and extraordinary waves and the gyro-frequency (cyclotron frequency) is also clearly noticeable and sounded in ionograms, in the following figure.

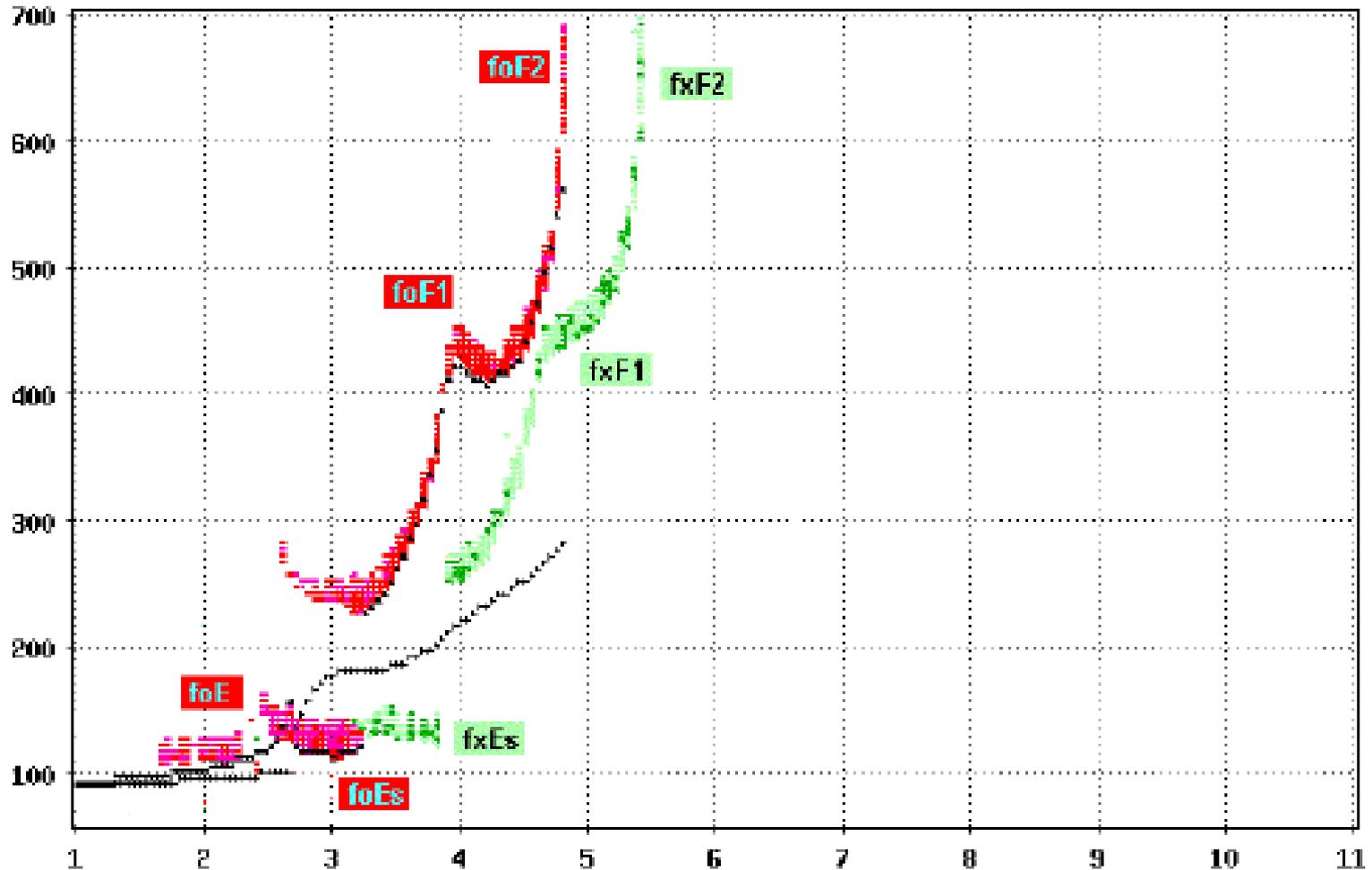


Ordinary & Extraordinary WAVES

STATION
Dourbes

YYYY DAY ODD HHMM P1 FFS S AMN PPS IGA PS
2004 Jun02 154 0600 MHH 400-1 8c5 100 +1+ R1

FoF2	4.60
FoF1	3.97
FoF1p	N/A
FoE	2.68
FoEp	2.60
FxT	5.50
FoEs	3.20
Fmin	1.60
<hr/>	
MUF	14.86
M	2.915
D	3000
<hr/>	
h'F	225
h'F2	485
h'E	90
h'Es	115
<hr/>	
znF2	262
znF1	217
znE	160
UF2	61
UF1	67
UE	14
B0	142.3
B1	1.40
<hr/>	
D-level	21



f_oF₂
f_oF₁
f_oF₁p
f_oE
f_oE_p
f_xT
f_oE_s
f_{min}
MUF
M
D
h'F
h'F₂
h'E
h'E_s
z_nF₂
z_nF₁
z_nE
U_F2
U_F1
U_E
B₀
B₁
D-level

Ordinary & Extraordinary WAVES

Both waves have different critical frequencies for the different respective ionospheric layers. For example, we find the **F2 layer**, respectively f_oE_2 and f_xF_2 . The difference between the gyro-frequency as,

$$f_H = 2(f_x - f_o)$$

OR

$$f_x - f_o = \frac{1}{2} f_H$$

The critical frequency of the extraordinary wave $(1/2) f_H$ higher than the critical frequency of the ordinary wave $[f_x = f_o + (1/2) f_H]$. It is obvious that both waves also have their own **MUF** (maximum usable frequency) values as well. The extraordinary wave MUF will always be higher.

$$MUF = \frac{f_c}{\sin \alpha}$$

Where, **alpha** is angle of incidence and f_c is the critical frequency, either f_o or f_x of the considered layer.

Ordinary & Extraordinary WAVES

Electromagnetic waves propagate through a **magneto-ionic medium** in the (so-called) free space mode as long as $n^2 > 0$; where n is the **refractive index**. When $n=0$ then **reflection occurs** (refracting downward). So we can use the **Appleton Equation** for a vertically propagating signal including the Earth's magnetic field and set the refractive index n to zero and solve for **X**.

Δ The Appleton - Hartree Equation

The Appleton – Hartree equation describes the **complex index of refraction** of the ionosphere in terms of the free motion of the electrons under the influence of thermal motion, geomagnetic fields and ionic collisions. Its successful application to the behavior of radio reflection demonstrated that free electrons in the ionosphere's F-layer in fact cause reflection.

$$n^2 = 1 - \frac{2X(1-X)}{2(1-X) - Y^2 \pm [Y^4 + 4(1-X)^2 Y^2]^{1/2}}$$

Ordinary & Extraordinary WAVES

The Appleton - Hartree Equation

$$n^2 = 1 - \frac{2X(1-X)}{2(1-X) - Y^2 \pm [Y^4 + 4(1-X)^2 Y^2]^{1/2}}$$

$X = 1$: The ordinary wave (o - wave)

$X = 1 - Y$: The extraordinary wave (x - wave)

$X = 1 + Y$: The **z - wave**

Where,

$$X = \frac{f_p}{f}, \quad Y = \frac{f_H}{f}$$

and

$$f_p = \left(\frac{Ne^2}{m\epsilon_0} \right)^{1/2}$$

$$f_H = \frac{eB}{2\pi m}$$

N - number densities of the electrons, e - charge of the electrons, m - mass of the electrons, ϵ_0 - free space permittivity, B - the Earth's Magnetic Field, f_p - plasma frequency, f_H - gyro-frequency.

The o and z modes have left hand and the x mode right hand **polarization with respect to the magnetic field direction**. Except at high latitudes the z mode is rarely observed in ionograms. The two prevailing modes o and x can be identified by the sense of rotation of the E-Field vector.

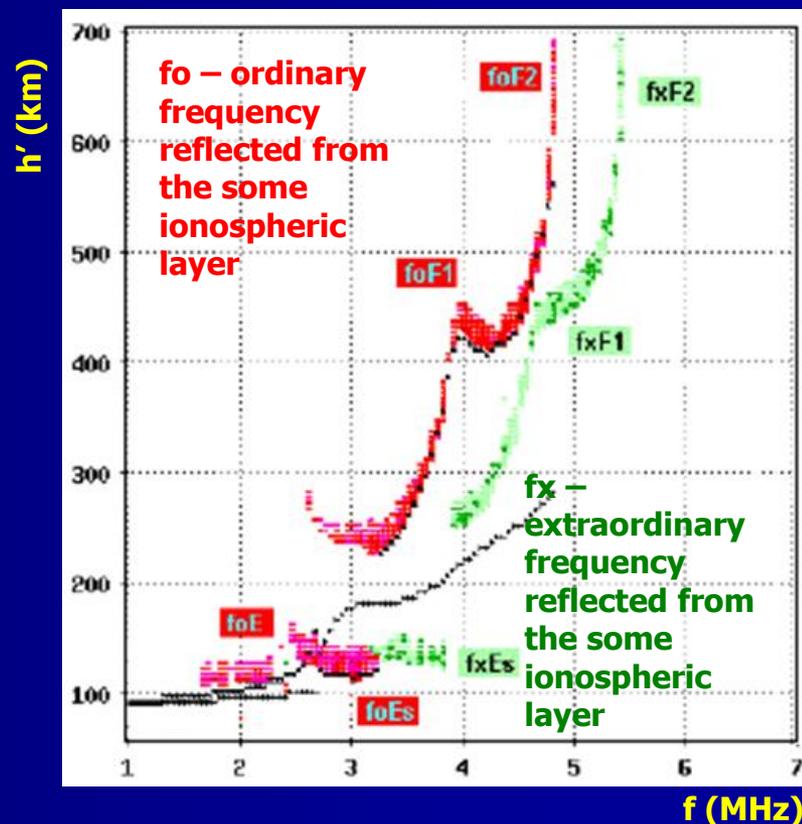
Ordinary & Extraordinary WAVES

Δ The deviation from great circle path

The illustration of the following figure help to display another important property of ordinary and extraordinary waves. The influence of the Earth's Magnetic Field causes them to split but also to divergence away from the great-circle path. The **ordinary wave** diverges toward the **magnetic pole** while the **extraordinary wave** diverges toward the (magnetic) **equator**.

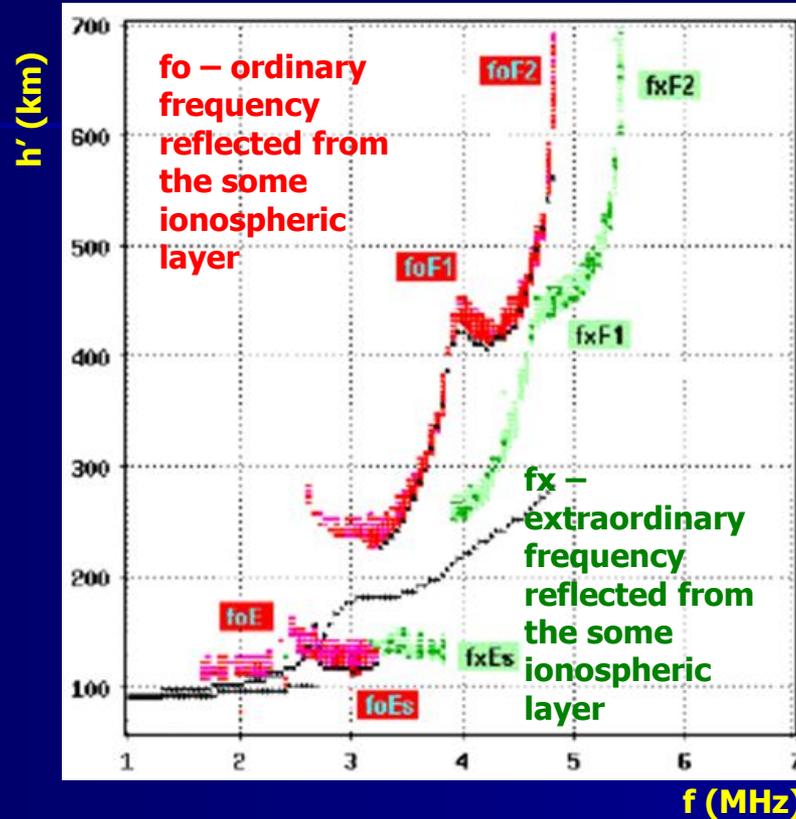
This can be observed in the ground track plots of these two signal components. NOTE: The divergence can also be cause additionally by ionospheric tilts (slope).

A typical ionogram showing the ordinary and extraordinary traces from the different ionospheric layers. This horizontal axis of the ionospheric sounder, and the vertical axis is the equivalent height.



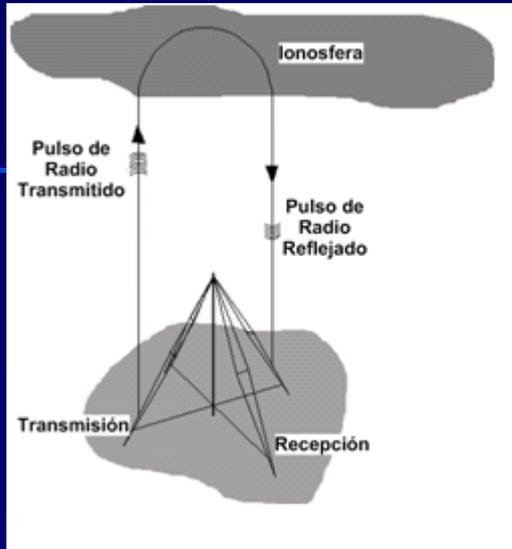
Ordinary & Extraordinary WAVES

Δ The deviation from great circle path



At frequencies lower than the peak plasma frequency (critical frequency) the ordinary wave deviates more pole-ward than the extraordinary wave equator-ward deflection (removal) does. The most deviation of the ordinary wave occurs also at or near the point of reflection.

Ionospheric Sounding Technique



In telecommunication and radio science, an ionospheric sounding is a technique that **provides real-time data** on high frequency ionospheric-dependent radio propagation, using a basic system consisting of a **synchronized transmitter** and a **receiver**.

The **time delay between transmission and reception is translated** into effective ionospheric layer altitude. Vertical incident sounding uses a collocated (arrange) transmitter and receiver and involving directing a range of frequencies vertically to the ionosphere and measuring the values of the reflected returned signals to determine the effective ionospheric layer altitude. This technique is also used to **determine the critical frequency**.

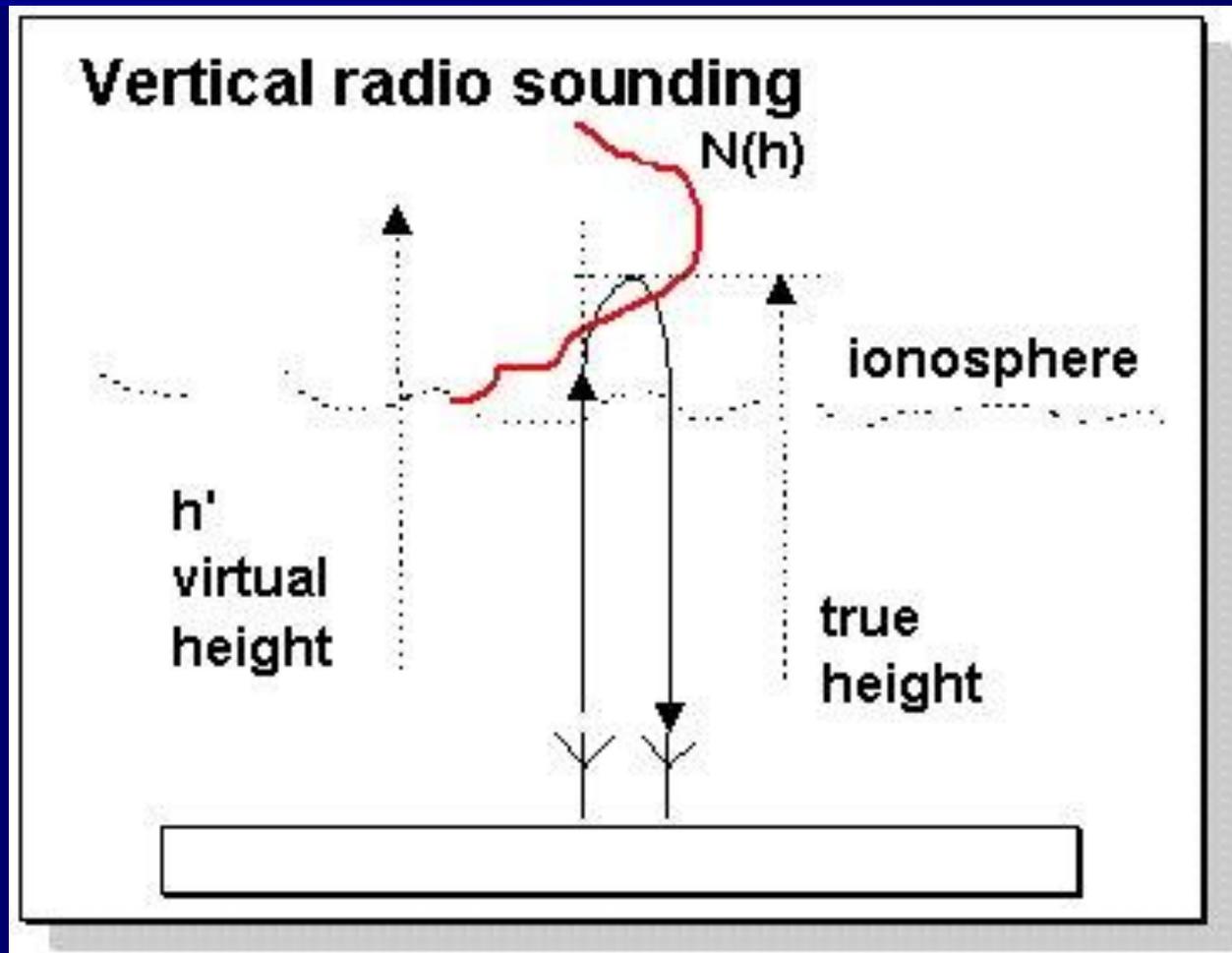
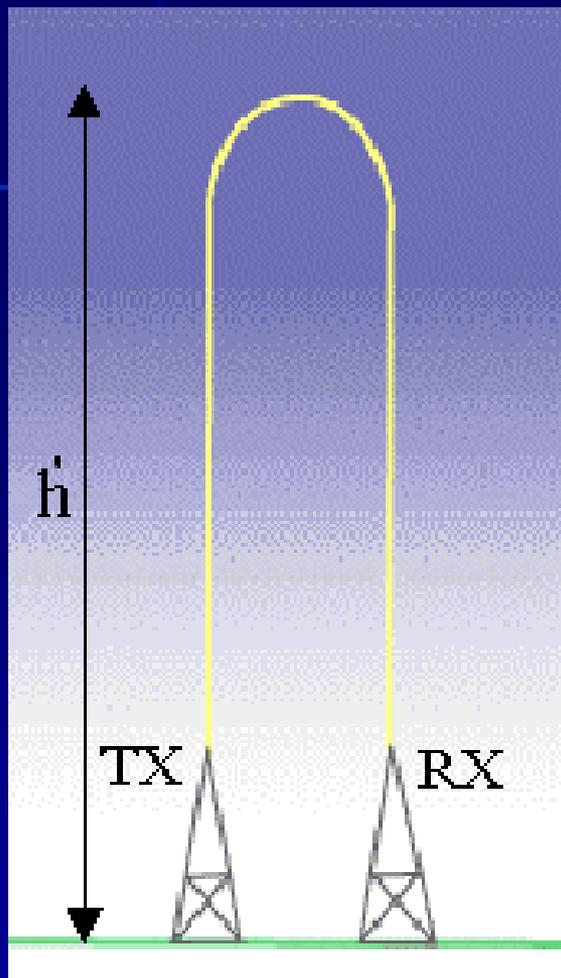
Ionospheric Sounding Technique

Oblique (indirect) sounders use a transmitter at one end of a given propagation path, and a synchronized receiver, usually with an oscilloscope type display (ionogram), at the other end.

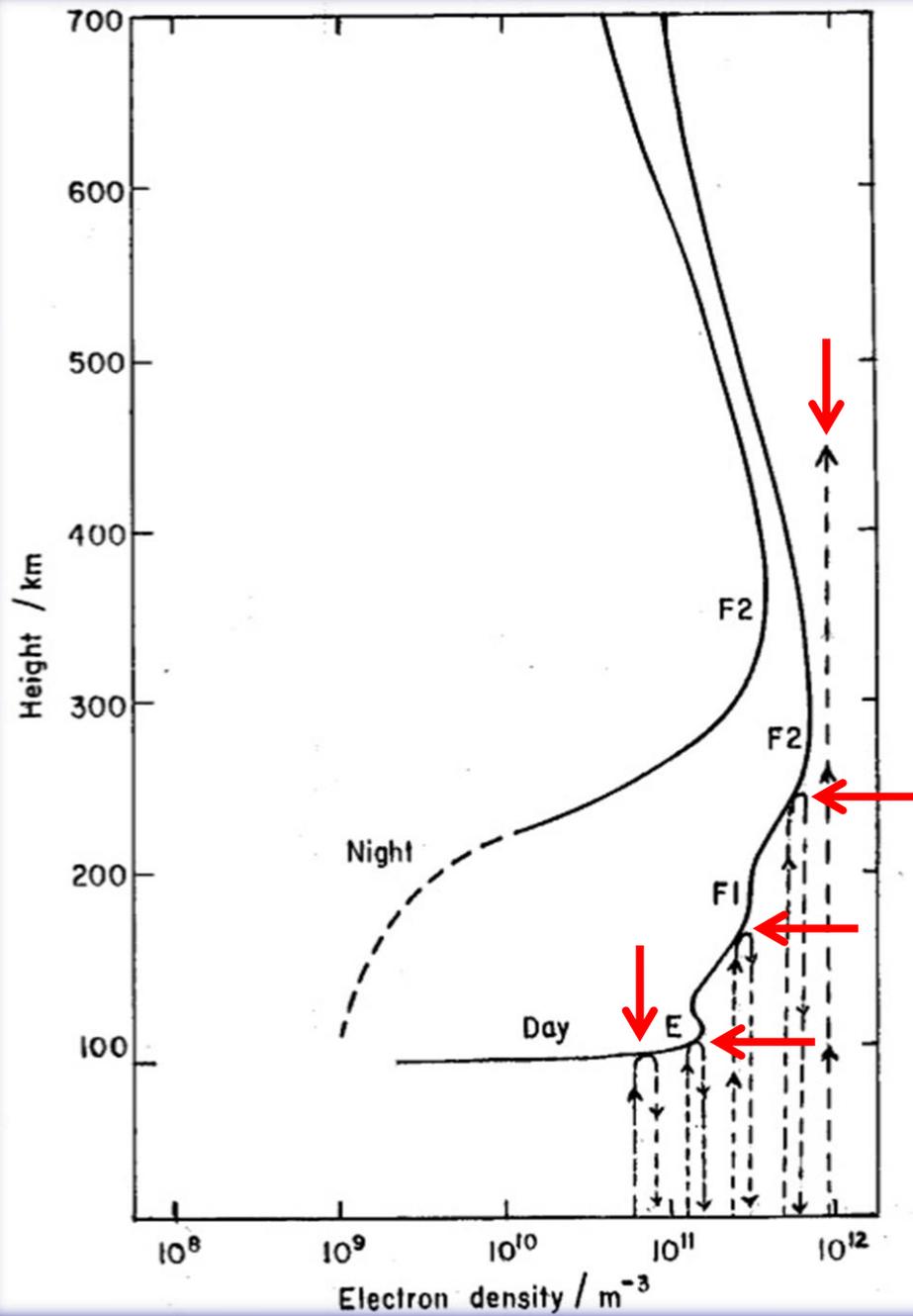


The transmitter emits a stepped (incident) or swept (sweep) frequency signal which is displayed or measured at the receiver. The measurement converts **time delay** to effective altitude of the ionospheric layer. The ionogram display shows the **effective altitude of the ionospheric layer** as a function of frequency.

Ionospheric Sounding Technique



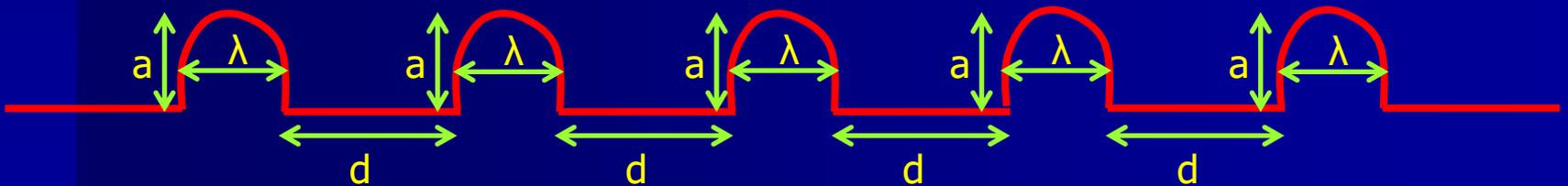
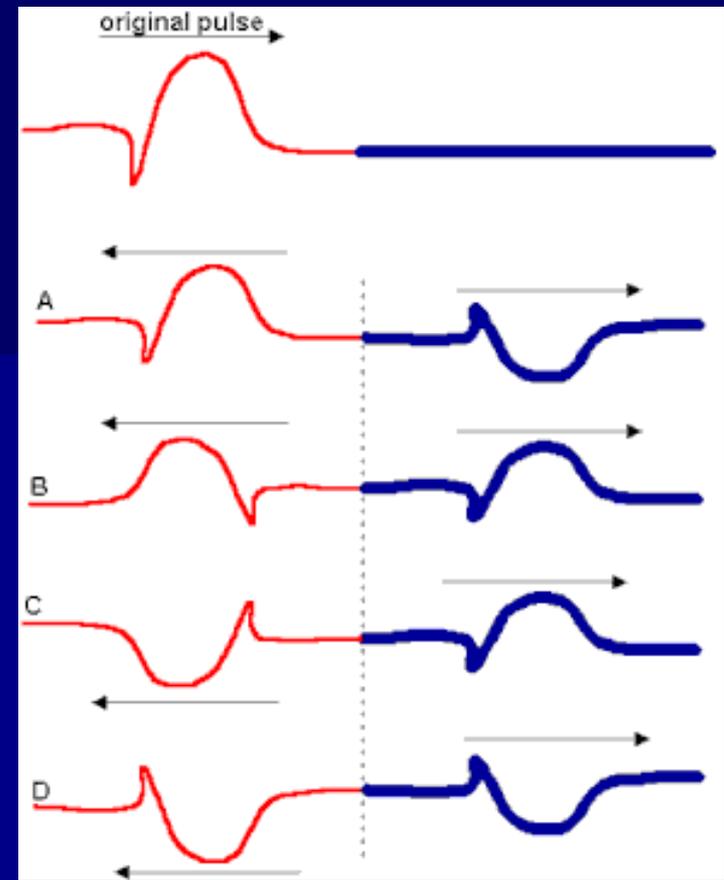
Ionospheric Sounding Technique



Pulse Reflection Method

Pulse reflection method is widely used to study the characteristics of the upper atmosphere by sending pulses towards it. These pulses are sequential and having the,

1. **same gaps between each pulses,**
2. **same amplitude**
- and
3. **same frequency.**



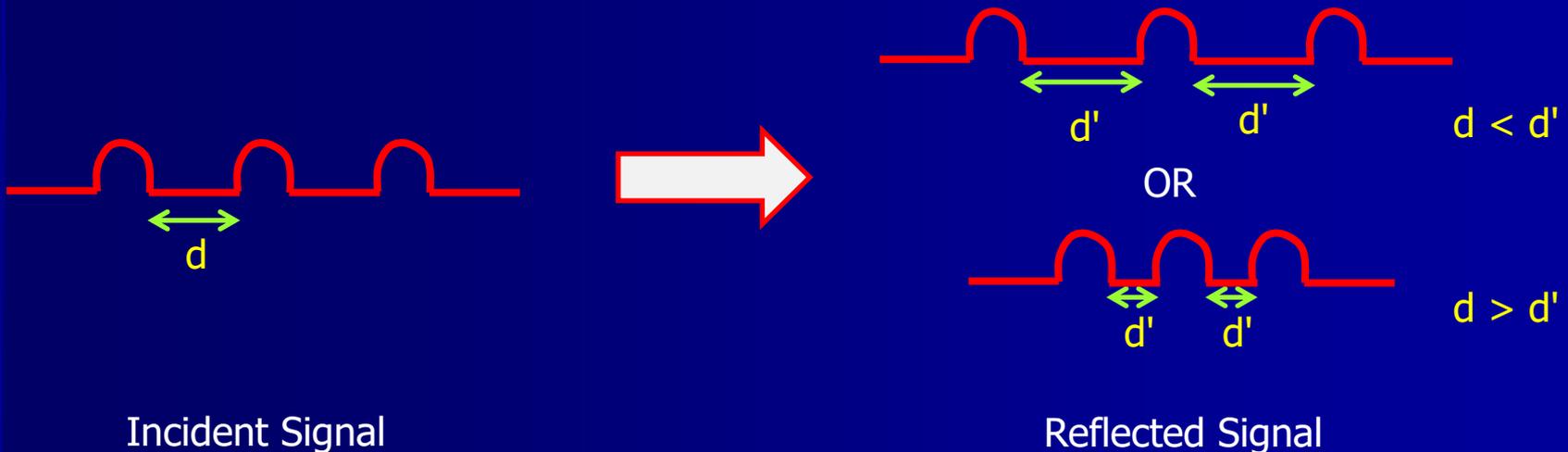
Where $c = f \lambda$

Pulse Reflection Method

Method :

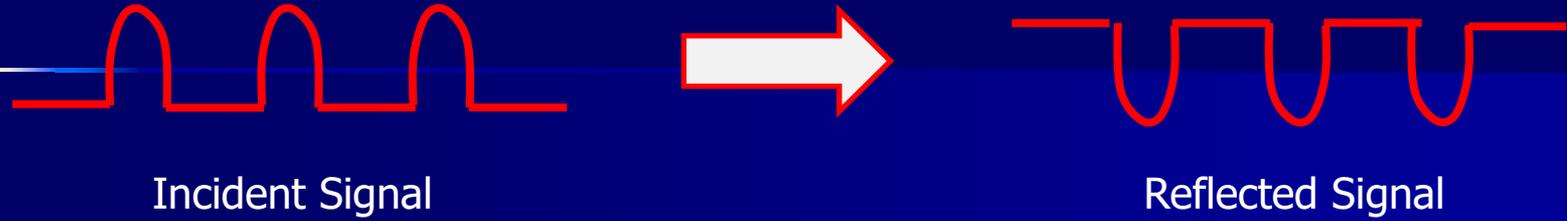
These pulses are sent towards the upper layers of the atmosphere and by observing the changes they show at the step of direction, then can be used to make conclusion regarding the characteristics of those layers. Such differences between the incident and reflected signals can be depicted as follows.

01. A difference can be observed between two consecutive gaps

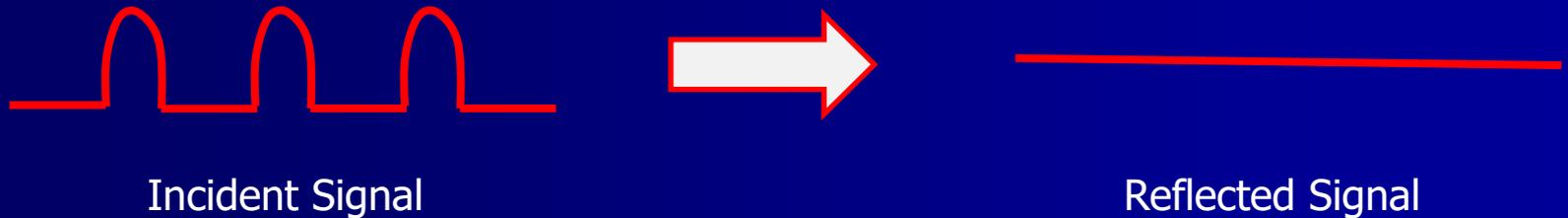


Pulse Reflection Method

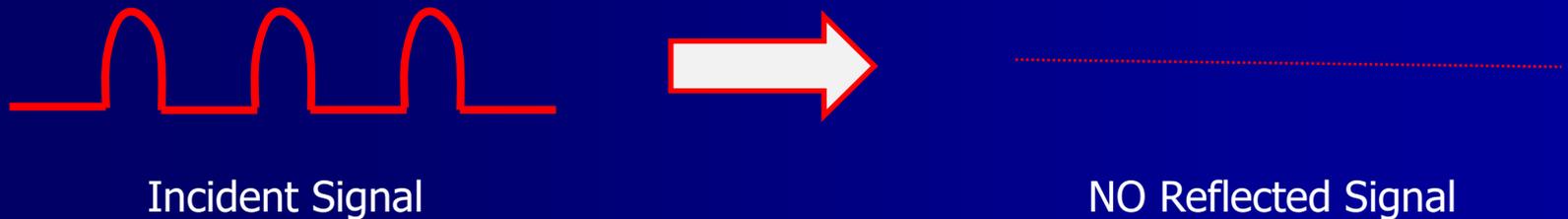
02. The wave form can be inverted during the step of reflection



03. The wave form can be vanished during the step of reflection

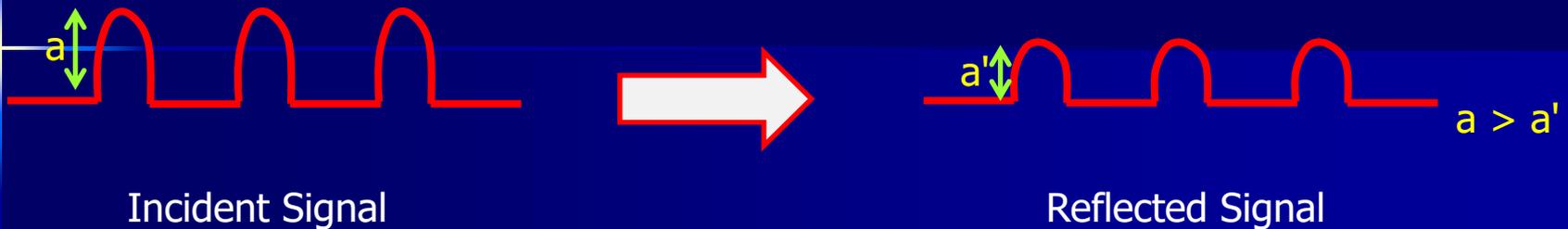


04. There is no reflected signal

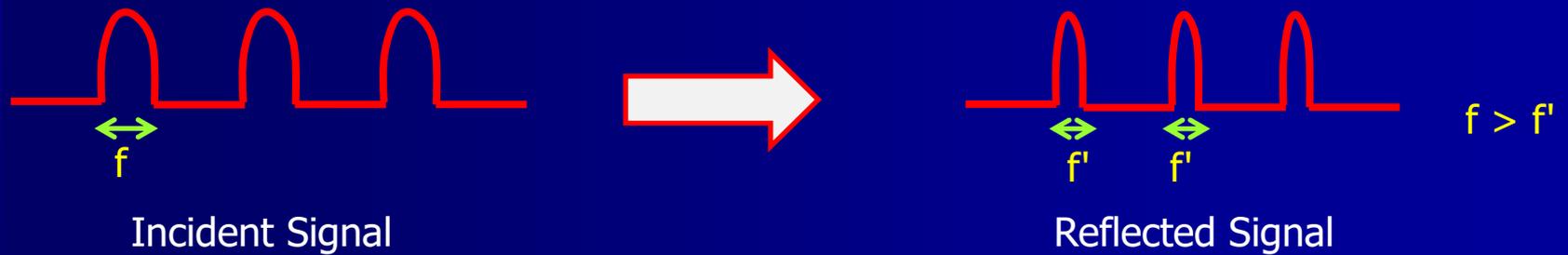


Pulse Reflection Method

05. Amplitude is tends to be reduced during the step of reflection



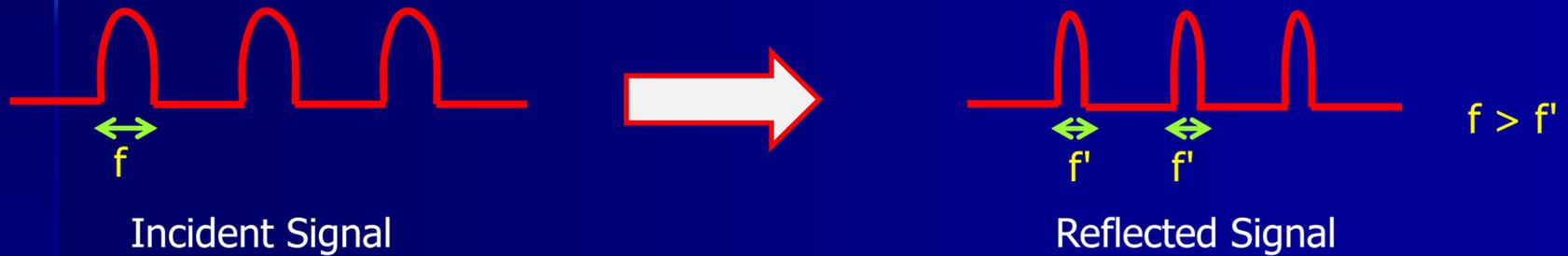
06. Frequency of the reflected pulse can be observed to be vary according to the prevailing circumstances in the upper atmosphere



Here the gap between two consecutive pulses wont get affected.

Pulse Reflection Method

06. Frequency of the reflected pulse can be observed to be vary according to the prevailing circumstances in the upper atmosphere



Here the gap between two consecutive pulses wont get affected.

Time that would take to detect after the reflection can be use to study the regions where these pulses being reflected from. **As the time vary**, it indicate that the **pulses are reflecting from different sections of the upper atmosphere !**

Thank You !

